

# Quantum Mechanics II

## WKB

**Derivation;  
Properties & Connection Formulae;  
Application**

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# WKB

## Derivation

*Extra!*

- 1926: Gregor **W**entzel, Hans **K**ramers & Léon **B**rillouin
- also: Francesco Carlini (1817), Joseph **L**iouville (1837), George **G**reen (1837), Lord Rayleigh (John Strutt, 1912), Richard Gans (1915), Harold **J**effreys (1923)
- Focus on 1-dimensional physics:

$$\hat{H}\psi = E\psi \quad \hat{H} = -\frac{\hbar^2}{2M} \frac{d^2}{dx^2} + W(x) \quad \psi(x) = e^{iu(x)}$$

- Substituting yields:

$$(u'(x))^2 = k^2(x) + iu''(x) \quad k(x) := \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$

- ...which is a non-linear, inhomogeneous 2nd order differential equation—*vs.* the linear, homogeneous Schrödinger equation.
- But, it admits an iterative (“fixed point”) method of solving.

# WKB

## Derivation

- Too hard to solve exactly. However, if  $k(x) = k$  were constant,
  - $u(x) = \pm k \cdot x + C$ , so  $u'(x) = \pm k$  and  $u''(x) = 0$ ,
  - and  $\psi(x) = A e^{ik \cdot x} + B e^{-ik \cdot x}$  would be the exact solution.
- For “sufficiently slowly-varying potentials,” start with

$$u'_0(x) = \pm k(x) \quad u_0(x) = \pm \int dx k(x) + C \quad u''_0(x) \mapsto 0 \text{ ignored!}$$

$$(u'(x))^2 = k^2(x) + iu''(x)$$

$$k(x) := \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$

**Extra!**

“piece-wise constant” potentials



# WKB

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and  $\psi(x) = A e^{ik \cdot x} + B e^{-ik \cdot x}$  would be the exact solution.

*“piece-wise constant” potentials*

For “sufficiently slowly-varying potentials,” start with

$$u'_0(x) = \pm k(x) \quad u_0(x) = \pm \int dx k(x) + C \quad u''_0(x) \mapsto 0 \text{ ignored!}$$

...and iterate

$$u''_0(x) = \pm k'(x) \quad u_1(x) = \pm \int dx \sqrt{k^2(x) + iu''_0(x)} + C$$

$$u_2(x) = \pm \int dx \sqrt{k^2(x) + iu''_1(x)} + C$$

...so the  $n^{\text{th}}$  iteration is

$$u_n(x) = \pm \int dx \sqrt{k^2(x) + iu''_{n-1}(x)} + C$$

If  $u_0(x)$  was a good start,  $u_\infty(x)$  is exact.

# WKB

$$u_1(x) = \pm \int dx \sqrt{k^2(x) + iu_0''(x)} + C$$

## Derivation

**Extra!**

- Still not soluble exactly, in general.
- Approximate:

$$\sqrt{k^2 + ik'} = k \sqrt{1 + \frac{ik'}{k^2}} \approx k \left[ 1 + \frac{ik'}{2k^2} + \dots \right] = k + \frac{i}{2} \frac{k'}{k} + \dots$$

$$\begin{aligned} u_1(x) &= \pm \int dx \sqrt{k^2 \pm ik'} \\ &\approx \pm \int dx k \left[ 1 \pm \frac{ik'}{2k^2} + \dots \right] \\ &= \pm \int dx k + \frac{i}{2} \int dx \frac{k'}{k} + \dots \\ &= \pm \int dx k + \frac{i}{2} \int \frac{dk}{k} + \dots \\ &= \pm \int dx k + \frac{i}{2} \ln [k(x)] + \dots \end{aligned}$$

$$\begin{aligned} \psi(x) &= e^{iu(x)} \\ &= e^{\pm i \int dx k - \frac{1}{2} \ln[k(x)] + C} \\ &= e^C e^{\pm i \int dx k} e^{-\frac{1}{2} \ln[k(x)]} \end{aligned}$$

$$\begin{aligned} \rightarrow & \frac{A}{\sqrt{k(x)}} e^{+i \int dx k} \\ & + \frac{B}{\sqrt{k(x)}} e^{-i \int dx k} \end{aligned}$$



# WKB

$$k(x) := \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$

$$\psi_{\text{WKB}}(x) = \frac{A}{\sqrt{k(x)}} e^{+i \int dx k(x)} + \frac{B}{\sqrt{k(x)}} e^{-i \int dx k(x)}$$

## Properties & Connection Formulae

**Extra!**

- Wherever  $E > W(x)$ , kinetic energy is positive ☺

- $k(x)$  is real and  $\psi_{\text{WKB}}(x)$  is oscillatory.

- Wherever  $E < W(x)$ , kinetic energy is negative ☹

- $k(x)$  is imaginary and  $\psi_{\text{WKB}}(x)$  is non-oscillatory.

- Write  $\kappa(x) := -ik(x) = \sqrt{\frac{2M}{\hbar^2} [W(x) - E]}$   $k(x) = i\kappa(x)$

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

- Transition from one to another:

- If  $W(x)$  goes across  $E$  discontinuously,  $W(x) \neq E$ ,  $\psi_{\text{WKB}}(x)$  is well-defined.

- If  $W(x)$  goes across  $E$  continuously,  $W(x_*) = E$ ,  $\psi_{\text{WKB}}(x_*) \rightarrow \infty$

$$[W(x) - E] = F_1(x - x_*) + \dots \quad \frac{1}{\sqrt{\sqrt{x - x_*}}} e^{\int dx \sqrt{x - x_*}} \xrightarrow{x \rightarrow x_*} \frac{1}{\sqrt[4]{0}} e^{0^{3/2}}$$

*not a pole; branch-cut*

# WKB

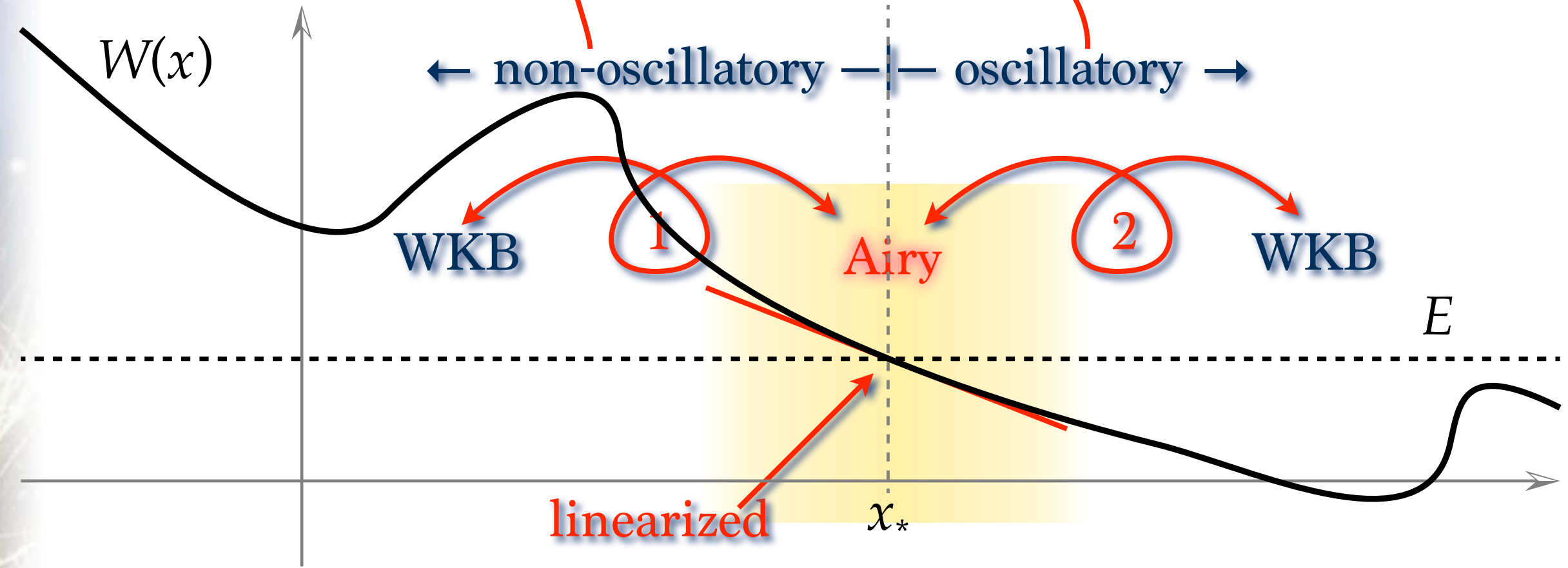
$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

$$\psi_{\text{WKB}}(x) = \frac{A}{\sqrt{k(x)}} e^{+i\int dx k(x)} + \frac{B}{\sqrt{k(x)}} e^{-i\int dx k(x)}$$

## Properties & Connection Formulae

*Extra!*

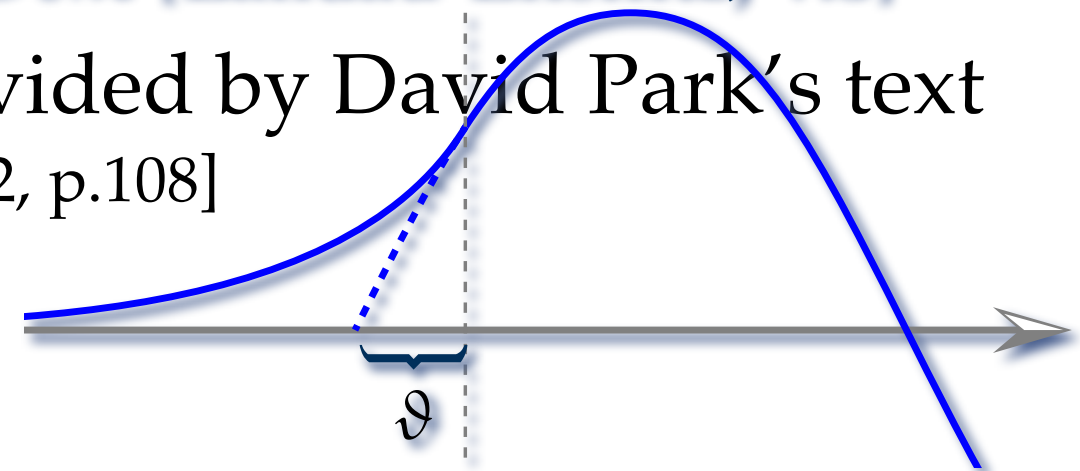
To match across  $x_*$



So, match:  $\text{WKB}_L$ —Airy— $\text{WKB}_R$  Done [Landau-Lifshitz, v.3]

Use the “connection formulae” provided by David Park’s text [Introduction to Quantum Physics, 3rd ed., 1992, p.108]

Phase-shift:  $\vartheta = e^{i\pi/4}$





# WKB

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

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## Properties & Connection Formulae

**Extra!**

For the non-oscillatory form:

- The C-term diverges for  $x \rightarrow -\infty$ ; the D-term diverges for  $x \rightarrow +\infty$
- All integrals are computed with  $x_*$  as the lower limit

Then:

### Barrier to Left

$$C = (\vartheta^* A + \vartheta B)$$

$$D = \frac{1}{2}(\vartheta A + \vartheta^* B)$$

$$A = \frac{1}{2}\vartheta C + \vartheta^* D$$

$$B = \frac{1}{2}\vartheta^* C + \vartheta D$$

### Barrier to Right

$$C = \frac{1}{2}(\vartheta^* A + \vartheta B)$$

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- The lower limit  $x^*$  in all integrals must refer to the transition point, *i.e.*, the classical turning point of the conn. formulae.
- In general: discontinuous  $W(x)/E$  crossing = no phase-shift
- ...continuous  $W(x)/E$  crossing =  $\pi/4$  phase-shift



# WKB

## Applications

### Smooth potential well

- Two transition points,  $a$  &  $b$
- Left transition, "Barrier to Left,"  $x_* = a$
- Right transition, "Barrier to Right,"  $x_* = b$

So,

$$\psi_L(x) = \frac{\cancel{C}}{\sqrt{\kappa(x)}} e^{-\int_a^x d\zeta \kappa(\zeta)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int_a^x d\zeta \kappa(\zeta)}$$

*diverges  $x < a$*

$$\psi_M(x) = \frac{A}{\sqrt{k(x)}} e^{+i\int_a^x d\zeta k(\zeta)} + \frac{B}{\sqrt{k(x)}} e^{-i\int_a^x d\zeta k(\zeta)}$$

$$= \frac{Ae^{+i\phi}}{\sqrt{k(x)}} e^{+i\int_b^x d\zeta k(\zeta)} + \frac{Be^{-i\phi}}{\sqrt{k(x)}} e^{-i\int_b^x d\zeta k(\zeta)}$$

$$\psi_R(x) = \frac{C'}{\sqrt{\kappa(x)}} e^{-\int_b^x d\zeta \kappa(\zeta)} + \frac{\cancel{D'}}{\sqrt{\kappa(x)}} e^{+\int_b^x d\zeta \kappa(\zeta)}$$

*diverges  $x > b$*

$A = A(D)$  &  $B = B(D)$ ,  $C' = C'(Ae^{+i\phi}, Be^{-i\phi})$  and  $D'(Ae^{+i\phi}, Be^{-i\phi}) = 0!$

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

$$\psi_{\text{WKB}}(x) = \frac{A}{\sqrt{k(x)}} e^{+i\int dx k(x)} + \frac{B}{\sqrt{k(x)}} e^{-i\int dx k(x)}$$

Barrier to Left

$$C = (\vartheta^* A + \vartheta B)$$

$$D = \frac{1}{2}(\vartheta A + \vartheta^* B)$$

$$A = \frac{1}{2}\vartheta C + \vartheta^* D$$

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$$\phi := \int_a^b d\zeta k(\zeta)$$

$$= \int_a^b d\zeta \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$

# WKB

## Applications

Left → Middle:  $A = \vartheta^* D$  &  $B = \vartheta D$

M → R:

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

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# WKB

## Applications

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$$C' = \frac{1}{2} (\vartheta^* (A e^{i\phi}) + \vartheta (B e^{-i\phi}))$$

$$= \frac{1}{2} (\vartheta^* (\vartheta^* D e^{i\phi}) + \vartheta (\vartheta D e^{-i\phi}))$$

$$= \frac{1}{2} (e^{-i\pi/2} D e^{i\phi} + e^{i\pi/2} D e^{-i\phi})$$

$$= D \frac{-ie^{i\phi} + ie^{-i\phi}}{2} = D \frac{e^{i\phi} - e^{-i\phi}}{2i} = D \sin(\phi)$$

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

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...and:  $D' = (\vartheta(A e^{i\phi}) + \vartheta^*(B e^{-i\phi}))$

$$= (\vartheta(\vartheta^* D e^{i\phi}) + \vartheta^*(\vartheta D e^{-i\phi}))$$

$$= (D e^{i\phi} + D e^{-i\phi})$$

$$= 2D \frac{e^{i\phi} + e^{-i\phi}}{2} = 2D \cos(\phi) \stackrel{!}{=} 0$$

# WKB

## Applications

$$\psi_{\text{WKB}}(x) = \frac{C}{\sqrt{\kappa(x)}} e^{-\int dx \kappa(x)} + \frac{D}{\sqrt{\kappa(x)}} e^{+\int dx \kappa(x)}$$

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...and:  $D' = (\vartheta(A e^{i\phi}) + \vartheta^*(B e^{-i\phi}))$

$$= (\vartheta(\vartheta^* D e^{i\phi}) + \vartheta^*(\vartheta D e^{-i\phi}))$$

$$= (D e^{i\phi} + D e^{-i\phi})$$

$$= 2D \frac{e^{i\phi} + e^{-i\phi}}{2} = 2D \cos(\phi) = 0$$

$$\phi = \begin{cases} (2n + \frac{1}{2})\pi \\ (2n + \frac{3}{2})\pi \end{cases}$$

$$\phi = \int_a^b dx \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$



# WKB

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$$\phi = \begin{cases} (2n + \frac{1}{2})\pi \\ (2n + \frac{3}{2})\pi \end{cases}$$

$$\phi = \int_a^b dx \sqrt{\frac{2M}{\hbar^2} [E - W(x)]}$$

Energy quantization

# WKB

## Applications

- Potential well w/ transition points  $x_* = a$  and  $x_* = b$
- The energy-quantization relation is:

$$\int_a^b dx \sqrt{\frac{2M}{\hbar^2} [E - W(x)]} = \begin{cases} (2n + \frac{1}{2})\pi \\ (2n + \frac{3}{2})\pi \end{cases} = \begin{cases} (2n + \frac{3}{4})\pi \\ (2n + \frac{7}{4})\pi \end{cases} = \begin{cases} (2n + 1)\pi \\ (2n + 2)\pi \end{cases}$$

symmetric  $\rightarrow$  (2n + 1/2)π, (2n + 3/4)π, (2n + 1)π  
 antisymmetric  $\rightarrow$  (2n + 3/2)π, (2n + 7/4)π, (2n + 2)π

- Whenever  $W(x)$  crosses  $E$  discontinuously

$$\Delta\psi(x) = 0 = \Delta\psi'(x)$$

- Whenever  $W(x)$  crosses  $E$  continuously

### WKB connection formulae

#### Barrier to Left

$$\begin{aligned} C &= (\vartheta^* A + \vartheta B) \\ D &= \frac{1}{2}(\vartheta A + \vartheta^* B) \\ A &= \frac{1}{2}\vartheta C + \vartheta^* D \\ B &= \frac{1}{2}\vartheta^* C + \vartheta D \end{aligned}$$

#### Barrier to Right

$$\begin{aligned} C &= \frac{1}{2}(\vartheta^* A + \vartheta B) \\ D &= (\vartheta A + \vartheta^* B) \\ A &= \vartheta C + \frac{1}{2}\vartheta^* D \\ B &= \vartheta^* C + \frac{1}{2}\vartheta D \end{aligned}$$



## Quantum Mechanics II

*Now, go forth and  
calculate!!!*

**Tristan Hübsch**

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