

Quantum Mechanics I

Quantum Mechanics

The Schrödinger Equation:

Piecewise Constant Potentials

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Pink Floyd: "Another Brick in the Wall (Pt. II)"

Piecewise Constant Potentials

REMINDER

- Consider the Schrödinger equation

$$i\hbar \frac{\partial}{\partial t} \Psi(\vec{r}, t) = H \Psi(\vec{r}, t) \quad H = -\frac{\hbar^2}{2M} \vec{\nabla}^2 + W(\vec{r})$$

- Use stationary states

$$\Psi(\vec{r}, t) = e^{-iEt/\hbar} \psi_E(\vec{r}) \quad \left[-\frac{\hbar^2}{2M} \vec{\nabla}^2 + W(\vec{r}) \right] \psi_E(\vec{r}) = E \psi_E(\vec{r})$$

- Reduce to 1-dimensional space

$$k(x) := \frac{\sqrt{2M[E - W(x)]}}{\hbar} \quad \psi_E''(x) + \underbrace{\frac{2M}{\hbar^2} [E - W(x)]}_{:= [k(x)]^2} \psi_E(x) = 0$$

- With the $W(x)$ constant in some region, so is $k(x)$:

$$\begin{aligned} \psi(x) &= A \sin(kx) + B \cos(kx) &= A' \sinh(\kappa x) + B' \cosh(\kappa x) \\ &= C_+ e^{+ikx} + C_- e^{-ikx} &= C_+ e^{+\kappa x} + C_- e^{-\kappa x} \\ &= N \sin(kx + \delta) \end{aligned}$$

$$\kappa := ik$$

Piecewise Constant Potentials

REMINDER & A 2ND LOOK

$$\psi_E''(x) + \frac{2M}{\hbar^2} [E - W(x)] \psi_E(x) = 0$$

- Orthogonality and normalizability:

$$\int_{\text{all space}} dx \psi_{E'}^*(x) \psi_E(x) = 0 \quad \text{if } E' \neq E \quad \text{if } E' = E, \text{ normalize}$$

- ...depending on the energy spectrum:

$$\int_{\text{all space}} dx \psi_{E'}^*(x) \psi_E(x) = \begin{cases} \delta_{E',E} & \text{if } E \text{ is discrete,} \\ \delta(E' - E) & \text{if } E \text{ is continuous.} \end{cases}$$

- Matching conditions

- Everywhere continuous: $\Delta\psi_E(x) := \lim_{\epsilon \rightarrow 0} [\psi_E(x+\epsilon) - \psi_E(x-\epsilon)] = 0$

- Unless $W(x)$ crosses E continuously*,

- Wherever $\Delta W(x) < \infty$, $\Delta\psi_E'(x) = 0$

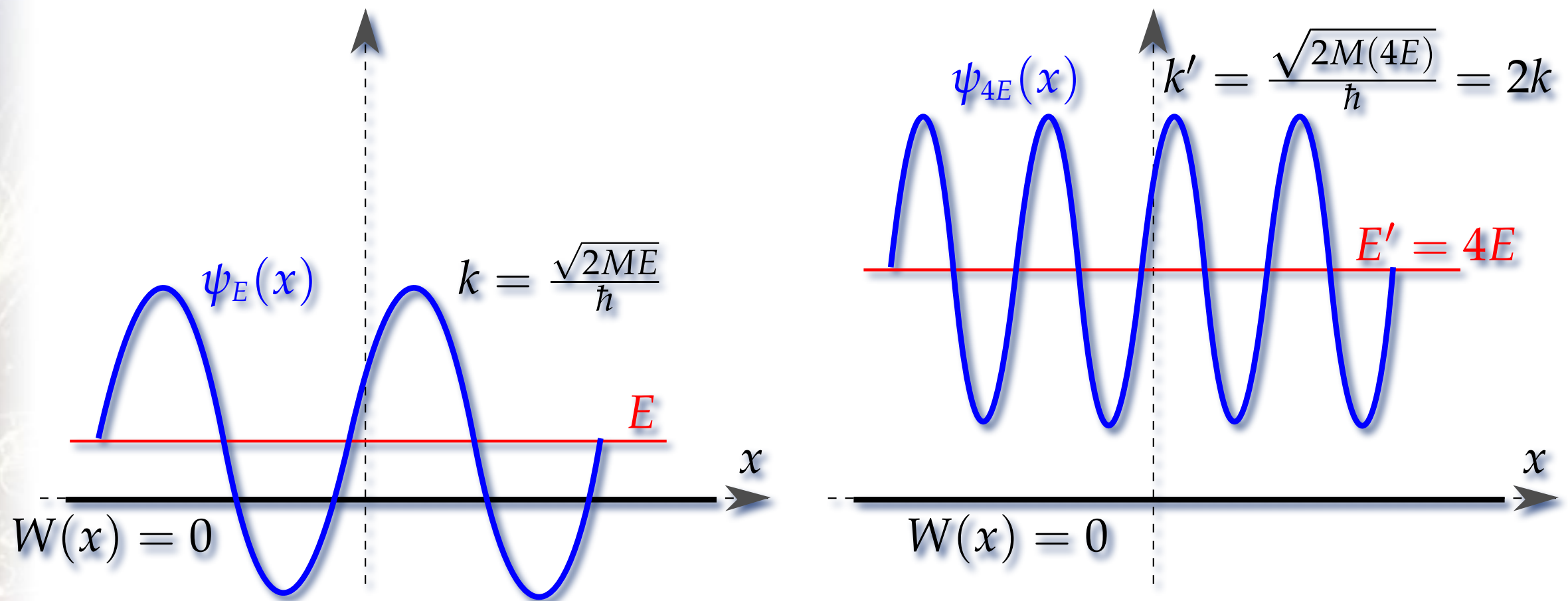
- Where $W(x) = w(x) + \lambda \delta(x-a)$, $\Delta\psi_E'(x) = \frac{2M\lambda}{\hbar^2} \psi_E(a)$

- Where $W(x) = \infty$, for $x \in [a, b]$, $\psi_E(x) = 0$, $\psi_E'(x)$ free

Piecewise Constant Potentials

A FEW POINTERS...

- Sketch: x horizontally, $W(x)$ vertically w/ E as a reference

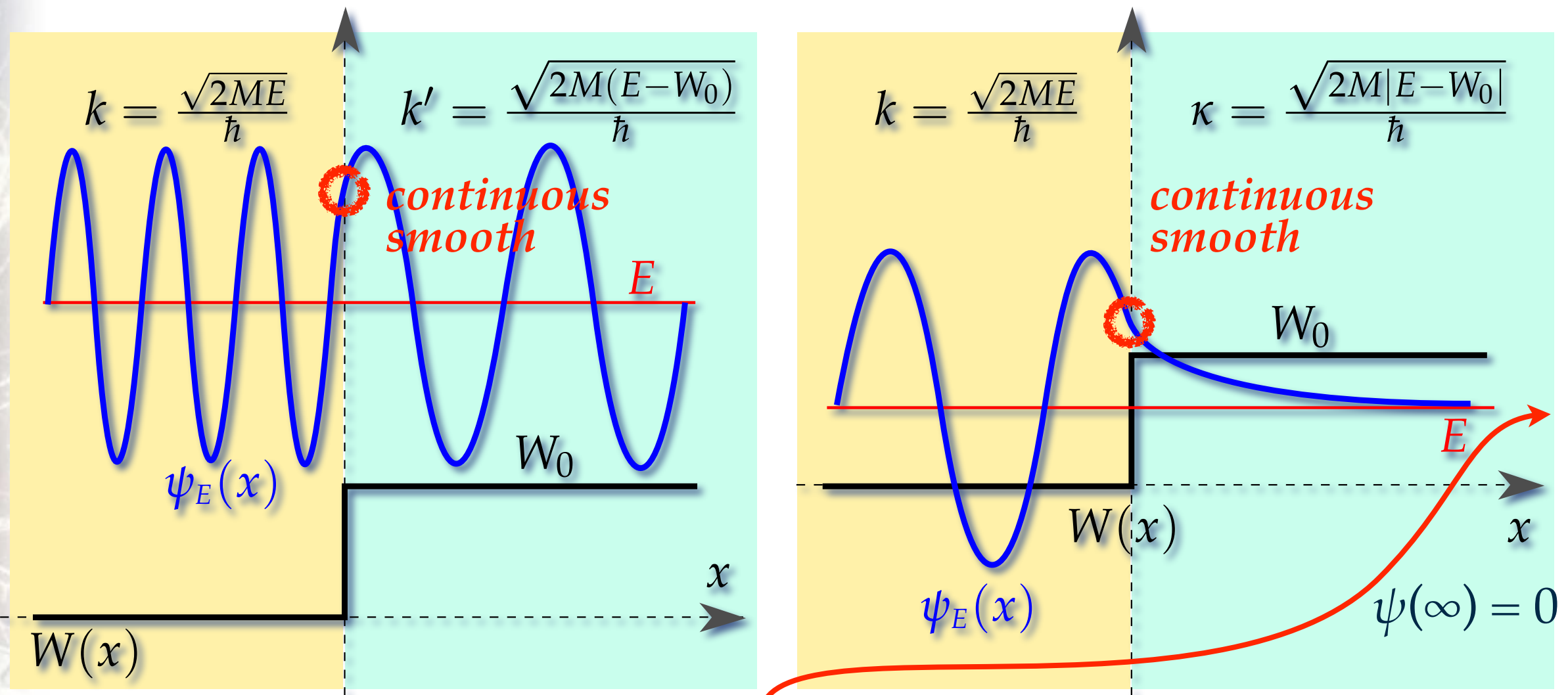


- Use the E -level as the horizontal axis for $\psi_E(x)$
- Solutions are oscillatory since $E - W(x) > 0$, and k is real
- The wave-number grows with the square-root of $E - W(x)$

Piecewise Constant Potentials

A POTENTIAL STEP: SKETCH

- The step function: $W(x)=0$ for $x<0$, $W(x)=W_0>0$ for $x>0$
- Region I: $x\in(-\infty,0)$, and region II: $x\in(0,+\infty)$



- Matching: @ $x=0$, and @ $x \rightarrow \infty$ only if $E < W_0$.

Piecewise Constant Potentials

A POTENTIAL STEP

- Now, do the math; $E > W_0$, using standard solutions:

$$\psi_I(x) = N \sin(kx + \delta)$$

$$\psi_{II}(x) = N' \sin(k'x + \delta')$$

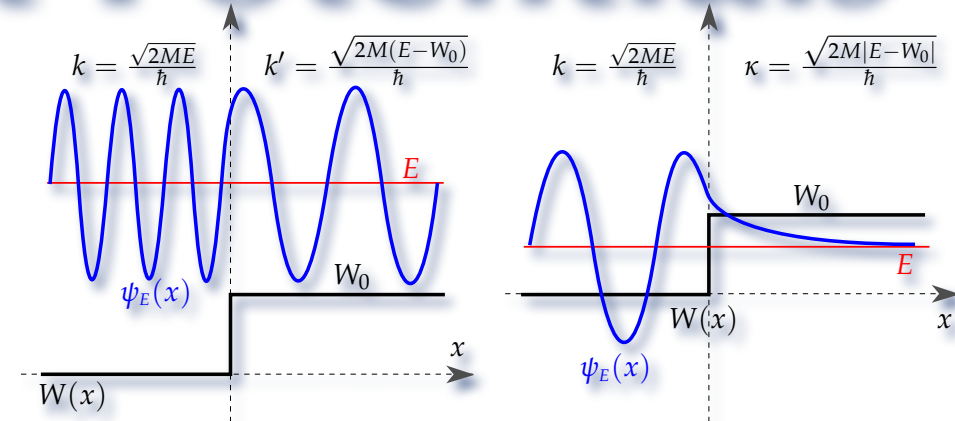
$$N \sin(\delta) = \psi_I(0) \stackrel{!}{=} \psi_{II}(0) = N' \sin(\delta') \quad \text{continuity}$$

$$kN \cos(\delta) = \psi'_I(0) \stackrel{!}{=} \psi'_{II}(0) = k'N' \cos(\delta') \quad \text{smoothness}$$

- Solve:

$$N' = N \frac{\sqrt{k^2 \cos^2(\delta) + k'^2 \sin^2(\delta)}}{k'} \quad \delta' = \tan^{-1} \left(\frac{k'}{k} \tan(\delta) \right) + 2n\pi$$

- Two conditions leave two integration constants, N and δ
- N is fixed by normalization; δ and E are free & continuous
- Continuous (unquantized!) energy is a hallmark of quantum scattering states (freely oscillate in all allowed directions)
- The dependence on E is convoluted, through k and k'



Piecewise Constant Potentials

A POTENTIAL STEP

- Now, do the math; $0 < E < W_0$, using standard solutions:

$$\psi_I(x) = N \sin(kx + \delta)$$

$$\psi_{II}(x) = \cancel{C_+} e^{+\kappa x} + C_- e^{-\kappa x}$$

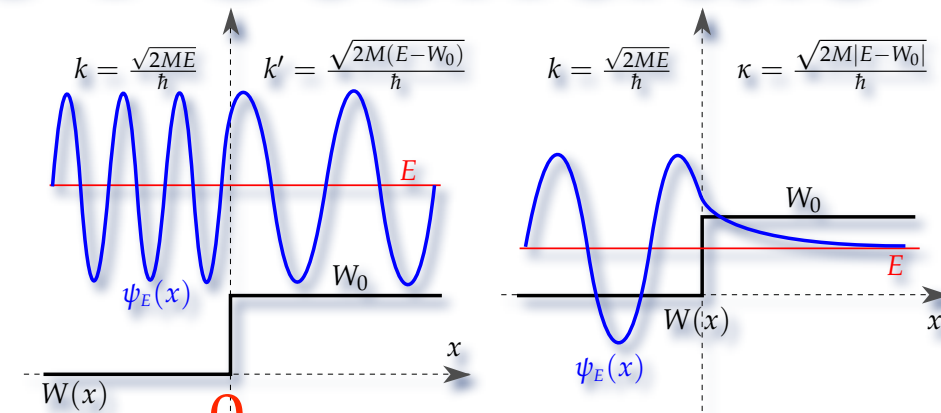
diverges keep

$$N \sin(\delta) = \psi_I(0) \stackrel{!}{=} \psi_{II}(0) = C_- \cdot 1$$

$$C_- = \frac{kN}{\sqrt{k^2 + \kappa^2}}$$

$$kN \cos(\delta) = \psi'_I(0) \stackrel{!}{=} \psi'_{II}(0) = -\kappa C_- \cdot 1$$

$$\delta = \tan^{-1}\left(\frac{k}{\kappa}\right) + 2n\pi$$

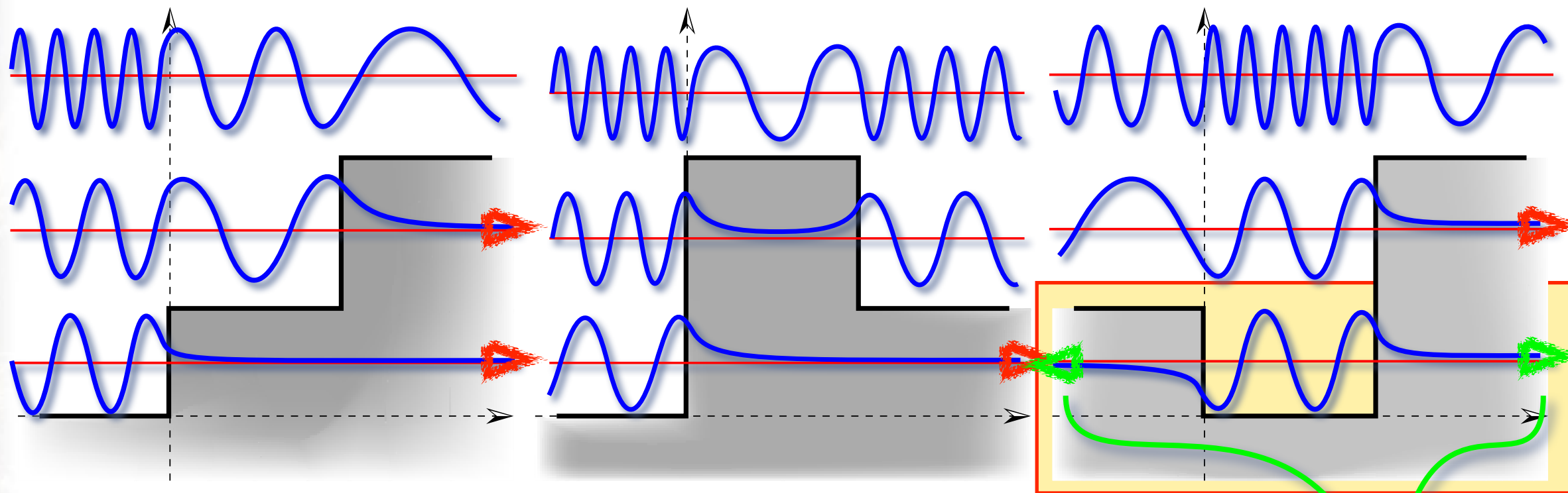


- Three conditions leave only one integration constant, N ,
- which is fixed by normalization; E is free & continuous
- though now the phase δ did become fixed
- Continuous (unquantized!) energy is a hallmark of quantum scattering states (freely oscillate in all allowed directions)
- The dependence on E is convoluted, through k and k'
- $1/\kappa$ is called the “skin depth”: the system penetrates the step

Piecewise Constant Potentials

TWO POTENTIAL STEPS: SKETCH

- Repeat the previous analysis in two steps
- Three radically different situations

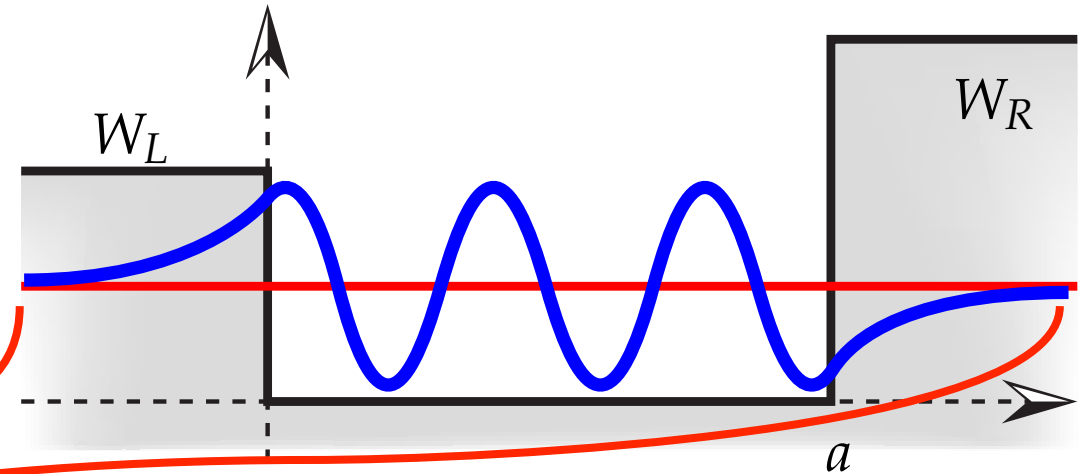


- Inter-region matching (cont's+smooth) leaves two constants
- Normalization fixes one constant (amplitude)
- Each **B.C. @ ∞ (normalizability)** fixes a constant } free E
- The lower-right case has "one too many conditions" } **fixes E**

Piecewise Constant Potentials

POTENTIAL WELL

- Do the math; $0 < E < \min(W_L, W_R)$ using standard solutions:



$$\psi_I(x) = C_+ e^{+\kappa x} + C_- e^{-\kappa x}$$

$$\psi_{II}(x) = A \cos(kx) + B \sin(kx)$$

$$\psi_{III}(x) = D_+ e^{+\kappa x} + D_- e^{-\kappa x}$$

$$\kappa = \sqrt{2M|E - W_L|/\hbar}$$

$$k = \sqrt{2ME/\hbar}$$

$$\kappa = \sqrt{2M|E - W_R|/\hbar}$$

- Matching:

$$0 \stackrel{!}{=} [A \cos(0) + B \sin(0)] - [C_+ e^0]$$

$$0 \stackrel{!}{=} [-kA \sin(0) + kB \cos(0)] - [\kappa C_+ e^0]$$

$$0 \stackrel{!}{=} [D_- e^{-\kappa a}] - [A \cos(ka) + B \sin(ka)]$$

$$0 \stackrel{!}{=} [-\kappa D_- e^{-\kappa a}] - [-kA \sin(ka) + kB \cos(ka)]$$

$$c := \cos(ka) \quad s := \sin(ka)$$

$$0 = \begin{bmatrix} -1 & 1 & 0 & 0 \\ -\kappa & 0 & k & 0 \\ 0 & -c & -s & e^{-\kappa a} \\ 0 & ks & -kc & -\kappa e^{-\kappa a} \end{bmatrix} \begin{bmatrix} C_+ \\ A \\ B \\ D_- \end{bmatrix}$$

det=0

A condition on E only

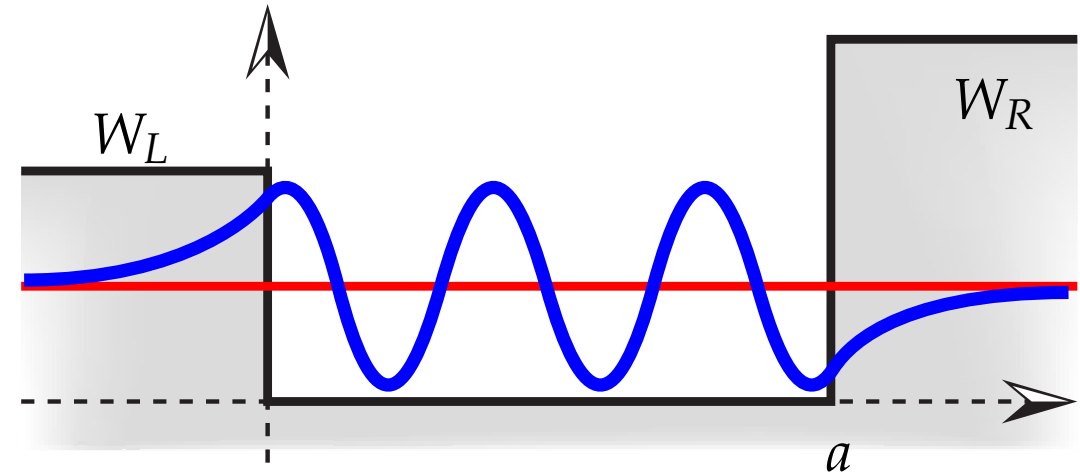
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POTENTIAL WELL

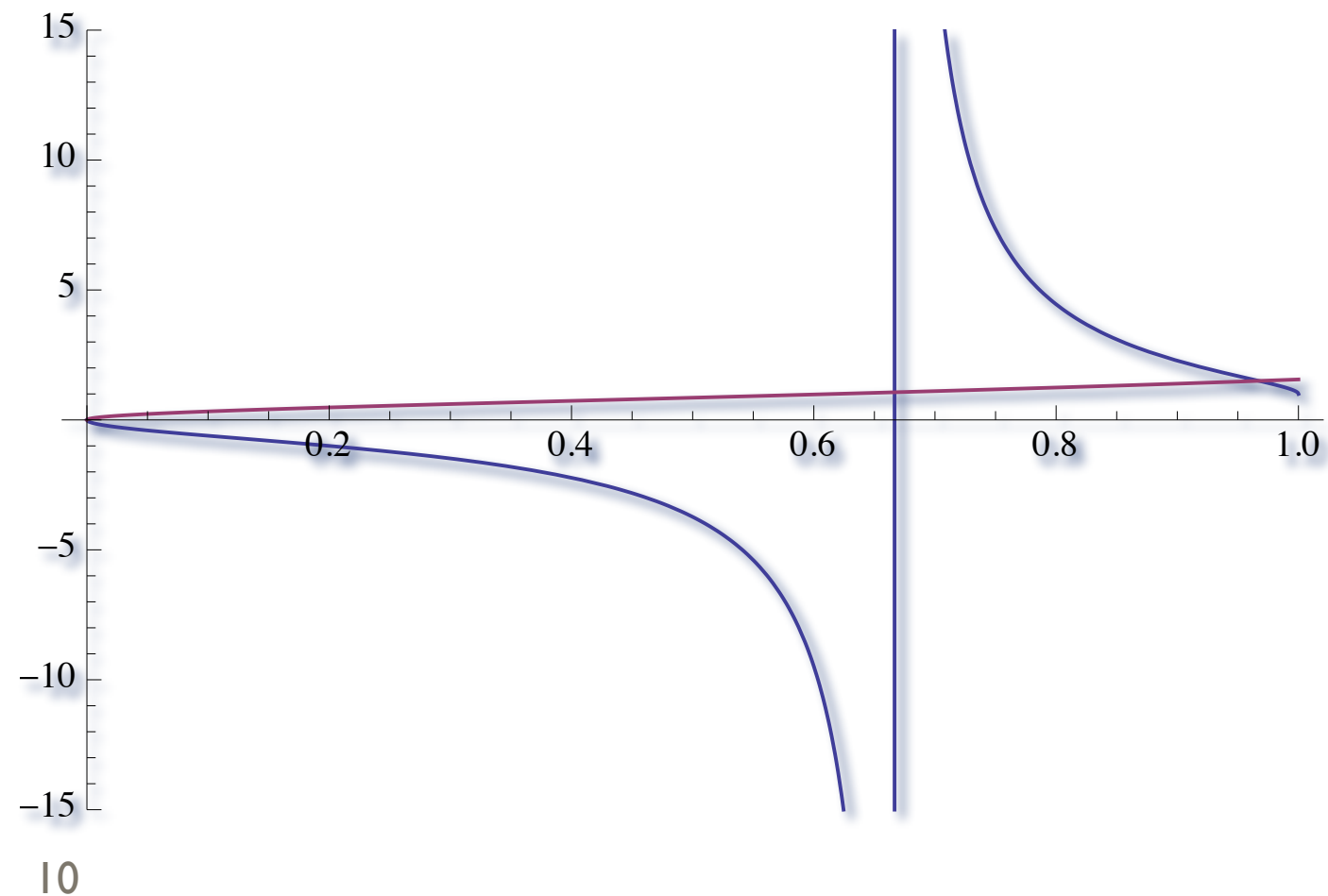
- The $\det=0$ condition gives

$$\tan(ka) = \frac{k(\kappa + \varkappa)}{k^2 - \kappa\varkappa}$$

$$\tan\left(\frac{\sqrt{2ME}a}{\hbar}\right) = \frac{\sqrt{ME}[\sqrt{M(W_L - E)} + \sqrt{M(W_R - E)}]}{ME - \sqrt{M(W_L - E)}\sqrt{M(W_R - E)}}$$



- This cannot be solved analytically
- ...but graphically + numerically



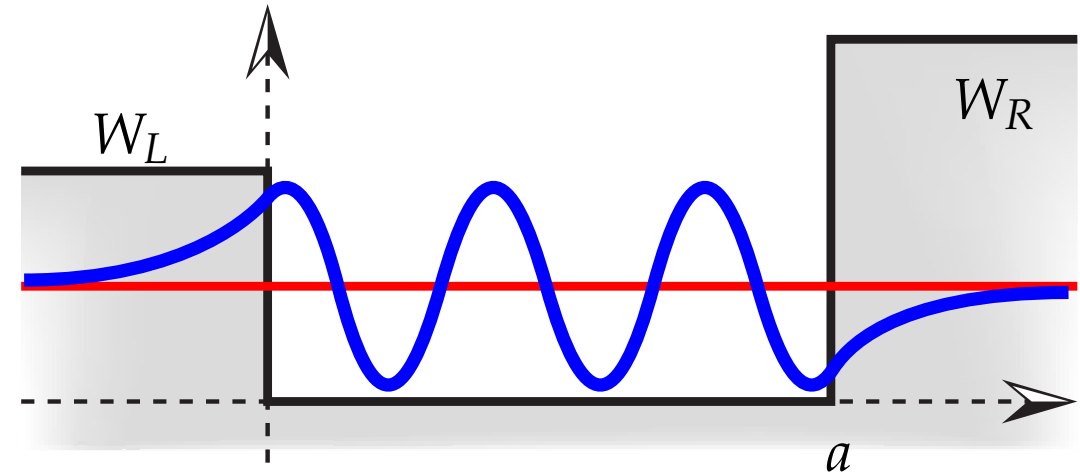
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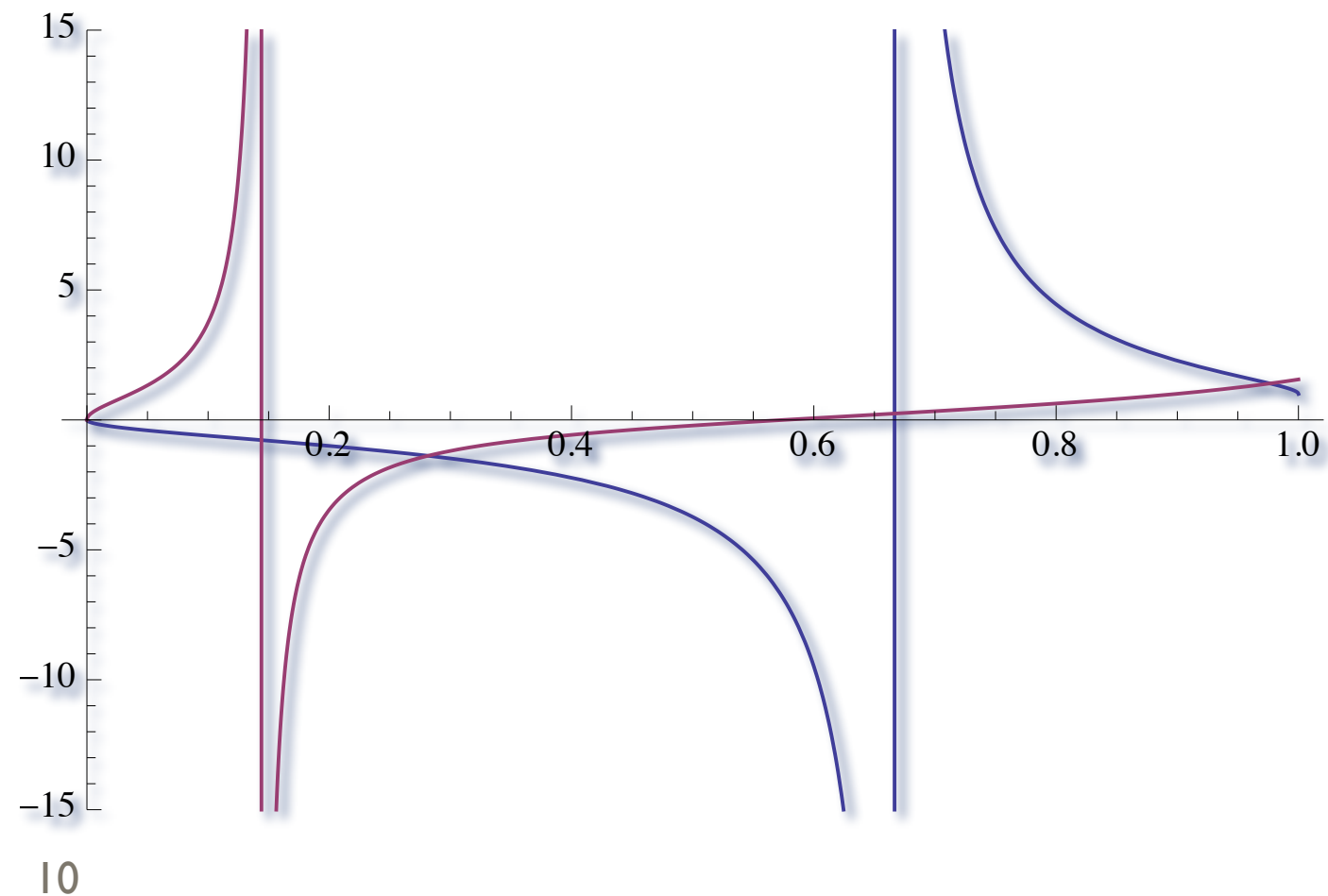
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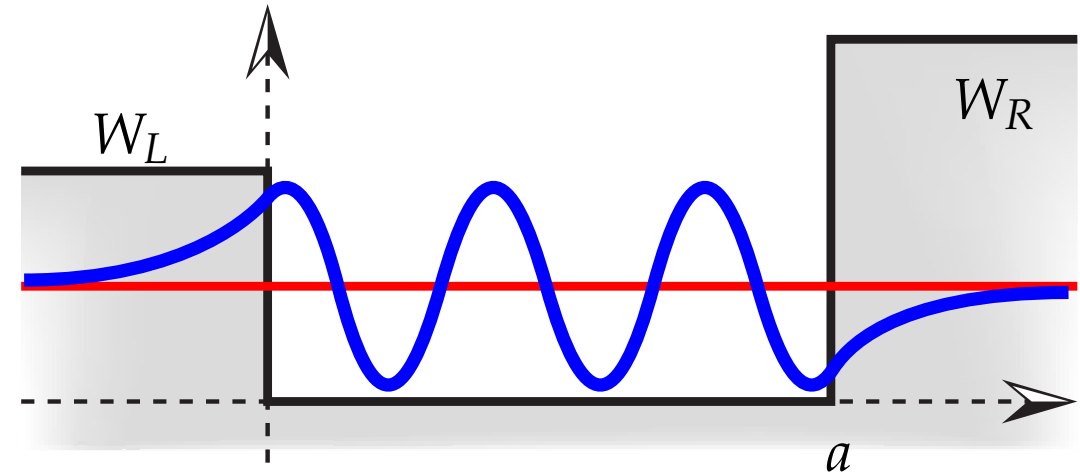
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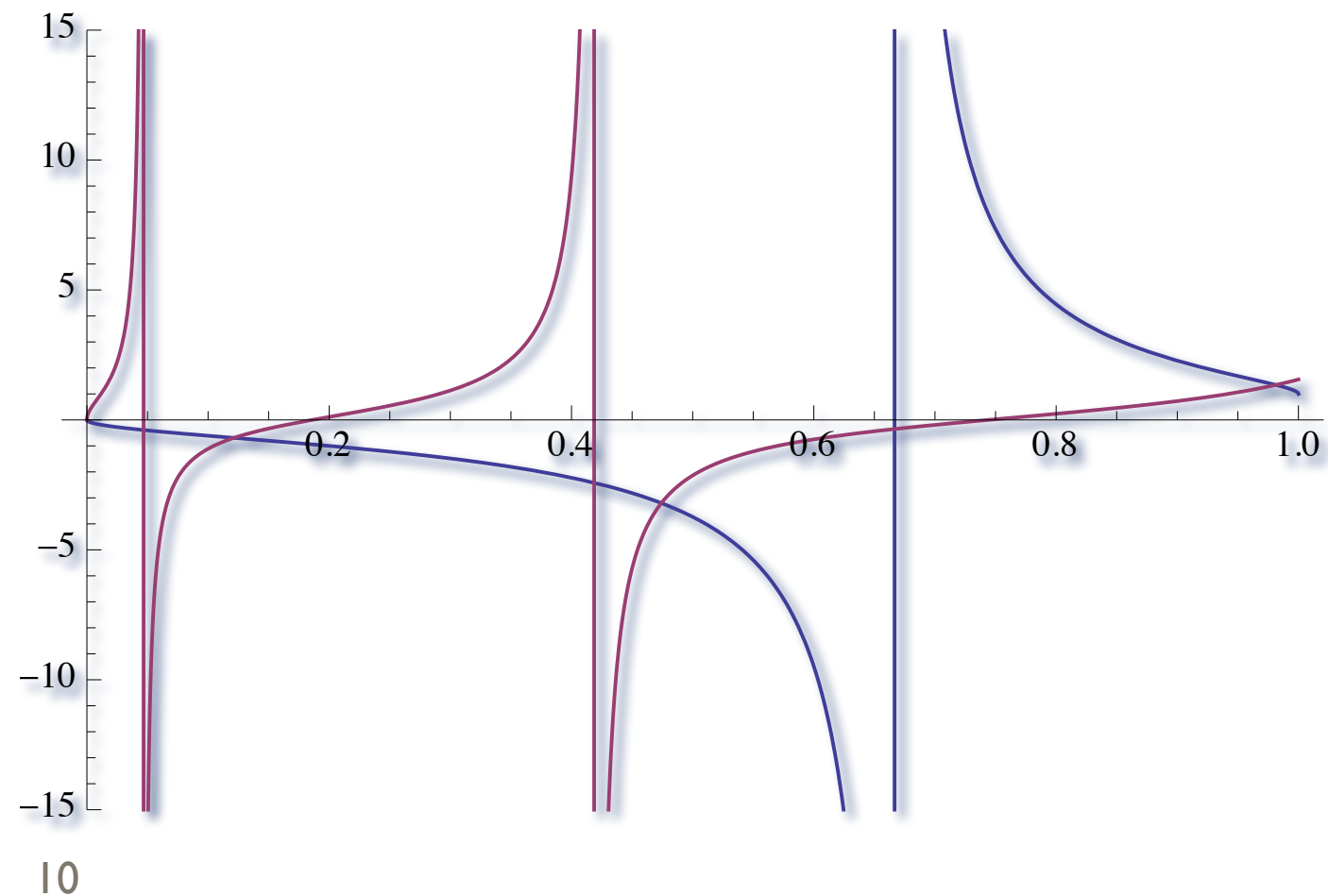
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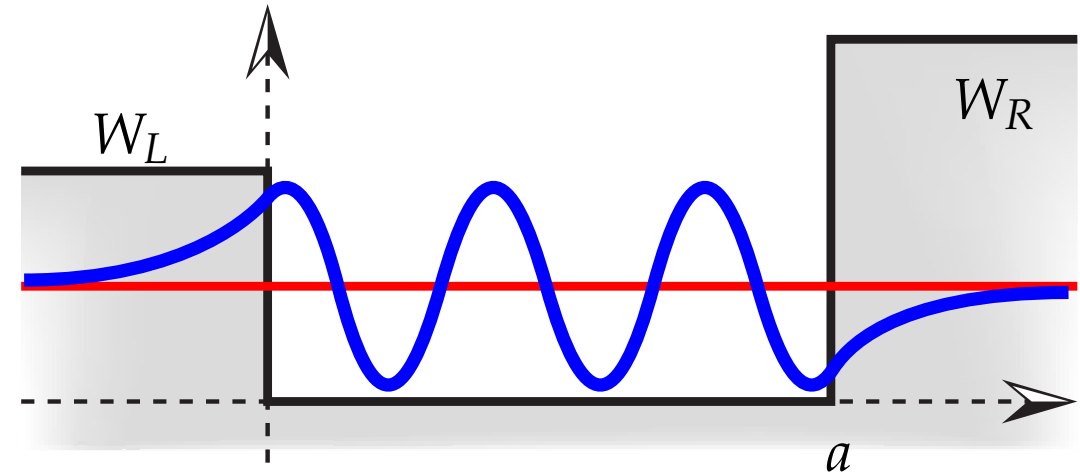
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POTENTIAL WELL

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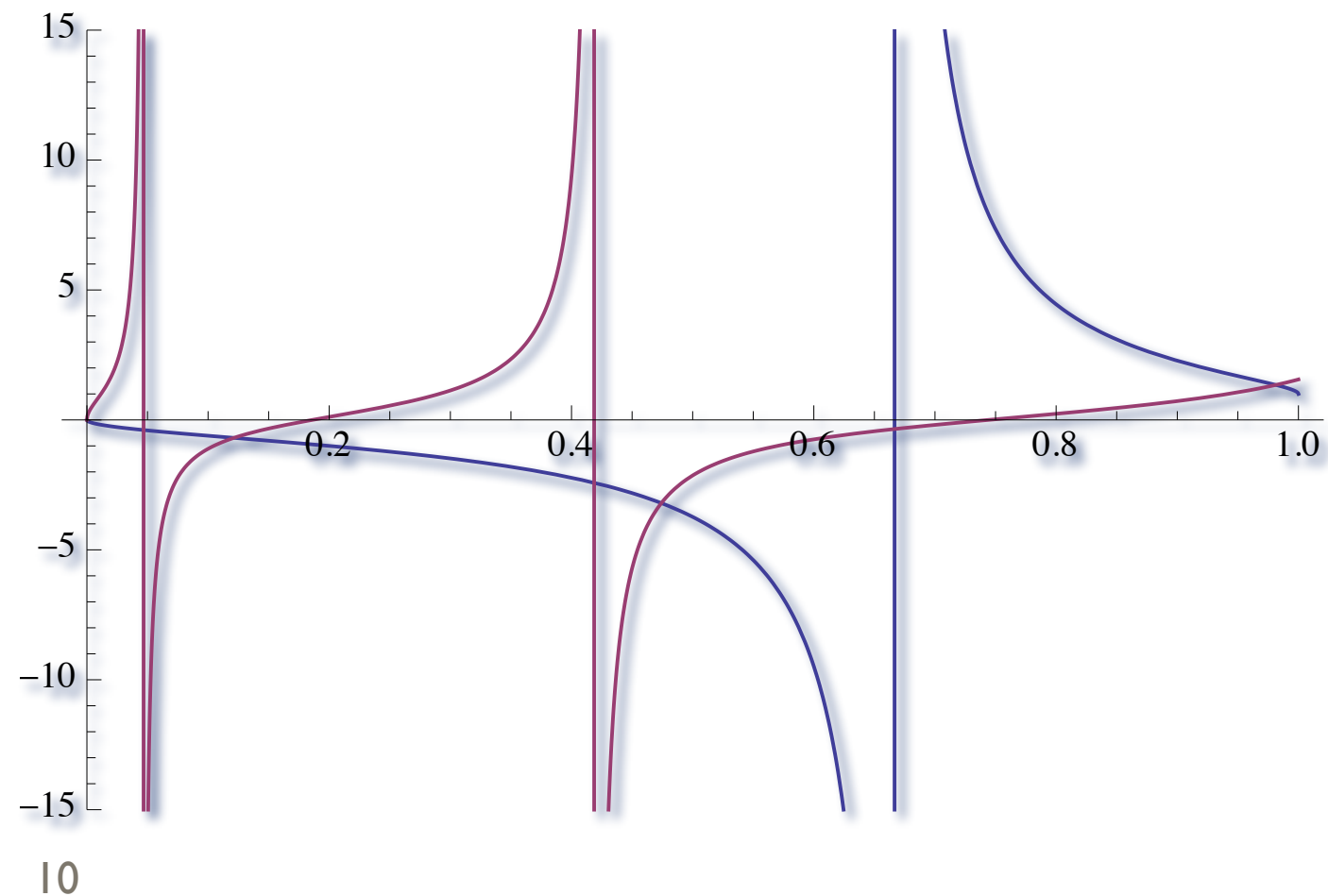
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- Lift $W_R \rightarrow \infty$:

$$\tan\left(\frac{\sqrt{2ME}a}{\hbar}\right) = -\sqrt{\frac{E}{W_L - E}}$$

- Lift also $W_L \rightarrow \infty$:

$$\tan\left(\frac{\sqrt{2ME}a}{\hbar}\right) = 0 \quad E_n = \frac{n^2 \pi^2 \hbar^2}{2M a^2}$$



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*Now, go forth and
calculate!!!*

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