

# (Fundamental) Physics of Elementary Particles

## Perturbation Theory in Pictures: Feynman Diagrams

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# Fundamental Physics of Elementary Particles

## PROGRAM

- Perturbation Theory in Pictures
  - Stationary perturbation theory in Quantum Mechanics
  - Feynman diagrams in perturbative Quantum Field Theory
    - The Heisenberg Zone — and What's Going On
    - A toy-model work-out
- A Few Stray Comments
  - An overview of the subject topic
  - A review of all conservation laws

# Perturbation Picturebook

## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Consider a quantum-mechanical system where the Hamiltonian differs from a well-known one by  $H'$ .
- We want to find solutions to:

$$H |n\rangle = E_n |n\rangle, \quad H := H_0 + \lambda H'$$

- ... where  $\lambda$  is a bookkeeping parameter.
- Given that we *do* know:

$$H_0 |n; 0\rangle = E_n^{(0)} |n; 0\rangle, \quad \begin{cases} \langle n; 0 | n'; 0 \rangle & = \delta_{n,n'}, \\ \sum_n |n; 0\rangle \langle n; 0| & = \mathbb{1}, \end{cases}$$

- ... we expand as a power-series in  $\lambda$ .

$$E_n = \sum_{k=0}^{\infty} \lambda^k E_n^{(k)}, \quad |n\rangle = \sum_{k=0}^{\infty} \lambda^k |n; k\rangle$$

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## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- We demand that order-by-order, the bases are orthonormalized (*always* doable, by Gram–Schmidt!)

$$\langle m; k | n; k \rangle = \delta_{mn}, \quad \forall m, n,$$

- but also that the “same” kets from different orders of perturbation are orthogonal:

$$\langle n; k | n; \ell \rangle = \delta_{kl}, \quad \forall k, \ell.$$

- Why? For consistency of the normalization!

$\langle n; k | n; k \rangle + 2 \Re(\lambda \langle n; k | n; k+1 \rangle)$   
 ~~$+ |\lambda|^2 \langle n; k+1 | n; k+1 \rangle$~~

$|n; k\rangle$   
 $|n; k+1\rangle$   
 $|n; k\rangle + \lambda |n; k+1\rangle$

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## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Now introduce specially weighted projectors:

$$\hat{\Pi}_n^\alpha := \sum_{m \neq n} \frac{|m; 0\rangle \langle m; 0|}{(E_n^{(0)} - E_m^{(0)})^\alpha}, \quad \text{so} \quad \hat{\Pi}_n^\alpha \hat{\Pi}_n^\beta = \hat{\Pi}_n^{\alpha+\beta}$$

- They project back to the well-known, initial basis, away from  $|n; 0\rangle$ , and curiously normalized.
- Note the no-degeneracy assumption!
- With this notation ( $\approx$  Cohen-Tannoudji, Diu, Laloë),

$$\left[ H_0 + \lambda H' \right] \left( \sum_{k=0}^{\infty} \lambda^k |n; k\rangle \right) = \left( \sum_{k'=0}^{\infty} \lambda^{k'} E_n^{(k')} \right) \left( \sum_{k''=0}^{\infty} \lambda^{k''} |n; k''\rangle \right)$$

- ... from where one obtains “order-by-order equations”:

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## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- The “order-by-order” equations:

$$|n; k\rangle = \hat{\Pi}_n^1 H' |n; k-1\rangle - \sum_{i=1}^{k-1} E_n^{(i)} \hat{\Pi}_n^1 |n; k-i\rangle, \quad k > 0,$$

$$E_n^{(k)} = \langle n; 0 | H' |n; k-1\rangle.$$

- The first few of these are easy:

$$E_n^{(1)} = \langle n; 0 | H' |n; 0\rangle,$$

$$|n; 1\rangle = \hat{\Pi}_n^1 H' |n; 0\rangle,$$

$$E_n^{(2)} = \langle n; 0 | H' \hat{\Pi}_n^1 H' |n; 0\rangle$$

$$|n; 2\rangle = \hat{\Pi}_n^1 (H' - E_n^{(1)}) |n; 1\rangle,$$

$$= \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n; 0\rangle - \hat{\Pi}_n^1 \hat{\Pi}_n^1 H' |n; 0\rangle \langle n; 0 | H' |n; 0\rangle,$$

$$= [\hat{\Pi}_n^1 H' \hat{\Pi}_n^1 - \hat{\Pi}_n^2 H' |n; 0\rangle \langle n; 0|] H' |n; 0\rangle,$$

... a subtraction!

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## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Soon enough, the complications grow:

$$\begin{aligned} E_n^{(3)} &= \langle n;0|H'|n;2\rangle \\ &= \langle n;0|H' [\hat{\Pi}_n^1 H' \hat{\Pi}_n^1 - \hat{\Pi}_n^2 H'|n;0\rangle \langle n;0|] H'|n;0\rangle \\ &= \langle n;0|H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H'|n;0\rangle - \langle n;0|H' \hat{\Pi}_n^2 H'|n;0\rangle \langle n;0|H'|n;0\rangle, \end{aligned}$$

- ... so the energy correction also requires subtractions ...

$$\begin{aligned} |n;3\rangle &= \hat{\Pi}_n^1 ((H' - E_n^{(1)})|n;2\rangle - E_n^{(2)}|n;1\rangle), \\ &= \hat{\Pi}_n^1 H'|n;2\rangle - \hat{\Pi}_n^1 |n;2\rangle \langle n;0|H'|n;0\rangle - \hat{\Pi}_n^1 |n;1\rangle \langle n;0|H' \hat{\Pi}_n^1 H'|n;0\rangle, \\ &= \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H'|n;0\rangle - \hat{\Pi}_n^1 H' \hat{\Pi}_n^2 H'|n;0\rangle \langle n;0|H'|n;0\rangle \\ &\quad - \hat{\Pi}_n^2 H' \hat{\Pi}_n^1 H'|n;0\rangle \langle n;0|H'|n;0\rangle - \hat{\Pi}_n^2 H'|n;0\rangle \langle n;0|H' \hat{\Pi}_n^1 H'|n;0\rangle \\ &\quad - \hat{\Pi}_n^3 H'|n;0\rangle \langle n;0|H'|n;0\rangle^2, \end{aligned}$$

- ... and the subtractions multiply combinatorially.
- The *combinatorial* growth gives a clue: **diagrams!**

# Perturbation Picturebook

## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Before we introduce any diagrams, however, consider:

$$\begin{aligned}
 |n;3\rangle &= \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle & |n;3\rangle &= \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^2 H' \hat{\Pi}_n^1 H' |n;0\rangle \langle n;0| H' |n;0\rangle & &- \hat{\Pi}_n^1 [H'] \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^2 H' |n;0\rangle \langle n;0| H' |n;0\rangle & &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 [H'] \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^2 H' |n;0\rangle \langle n;0| H' \hat{\Pi}_n^1 H' |n;0\rangle & &- \hat{\Pi}_n^1 [H' \hat{\Pi}_n^1 H'] \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^3 H' |n;0\rangle \langle n;0| H' |n;0\rangle^2, & &- \hat{\Pi}_n^1 [H'] \hat{\Pi}_n^1 [H'] \hat{\Pi}_n^1 H' |n;0\rangle \\
 & & & \left\{ \begin{aligned} &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 [H'] |n;0\rangle \\ &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 [H' \hat{\Pi}_n^1 H'] |n;0\rangle \end{aligned} \right.
 \end{aligned}$$

0



# Perturbation Picturebook

## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Before we introduce any diagrams, however, consider:

$$\begin{aligned}
 |n;3\rangle &= \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 \underline{[H']} \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 \underline{[H']} \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 \underline{[H']} |n;0\rangle \\
 &- \hat{\Pi}_n^1 \underline{[H']} \hat{\Pi}_n^1 \underline{[H']} \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 \underline{[H' \hat{\Pi}_n^1 H']} \hat{\Pi}_n^1 H' |n;0\rangle \\
 &- \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 \underline{[H' \hat{\Pi}_n^1 H']} |n;0\rangle
 \end{aligned}$$

An observation,  
not a proof!

The subtractions  
may in fact be  
generated from  
the original  
expression!

$$\hat{\Pi}_n^1 \underline{[H']} \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle = \hat{\Pi}_n^1 \cdot \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H' |n;0\rangle \langle n;0|H'|n;0\rangle,$$

$$\hat{\Pi}_n^1 \underline{[H' \hat{\Pi}_n^1 H']} \hat{\Pi}_n^1 H' |n;0\rangle = \hat{\Pi}_n^1 \cdot \hat{\Pi}_n^1 H' |n;0\rangle \langle n;0|H' \hat{\Pi}_n^1 H' |n;0\rangle,$$

- ... and so on. Now, you try for the 3<sup>rd</sup> one ...

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## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Recall:

$$\hat{\Pi}_n^\alpha |n;0\rangle = \sum_{m \neq n} \frac{|m;0\rangle \langle m;0|}{(E_n^{(0)} - E_m^{(0)})^\alpha} |n;0\rangle = \sum_{m \neq n} \frac{1}{(E_n^{(0)} - E_m^{(0)})^\alpha} |m;0\rangle \underbrace{\langle m;0|n;0\rangle}_{=0 (\because m \neq n)}.$$

$$|n;k\rangle = (\hat{\Pi}_n^1 H')^k |n;0\rangle - \text{all "excisions"}, \quad k \geq 0,$$

$$E_n^{(k)} = \langle n;0|H'(\hat{\Pi}_n^1 H')^{k-1}|n;0\rangle - \text{all "excisions"}. \quad k \geq 1,$$

- Now that we have the computations, depict them:

$$\begin{array}{c} \blacktriangleright \\ (in) \end{array} = |n;0\rangle, \quad \begin{array}{c} \blacktriangleright \\ (out) \end{array} = \langle n;0|,$$

$$\begin{array}{c} \blacktriangleright \\ (propagator) \end{array} = \hat{\Pi}_n^1, \quad \begin{array}{c} \blacktriangleright\blacktriangleright \\ (2^{nd} \text{ order propagator}) \end{array} = \hat{\Pi}_n^2, \quad \begin{array}{c} \otimes \\ (interaction) \end{array} = H'.$$

# Perturbation Picturebook

## QUANTUM MECHANICS: WARM-UP CALISTHENICS

- Soo...

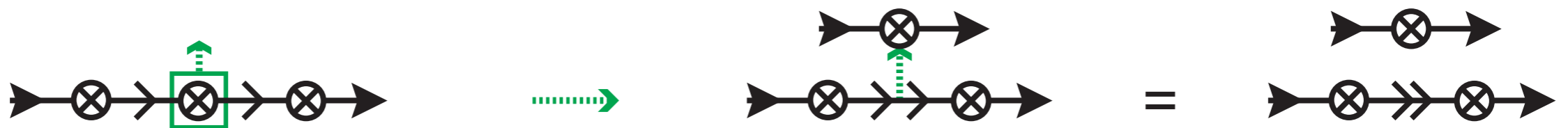
$$\langle n;0|H'|n;0\rangle \xrightarrow{(1.89a)} E_n^{(1)} = \begin{array}{c} n \quad n \\ \longrightarrow \otimes \longrightarrow \end{array},$$

$$\hat{\Pi}_n^1 H'|n;0\rangle \xrightarrow{(1.89b)} |n;1\rangle = \begin{array}{c} n \quad m \\ \longrightarrow \otimes \longrightarrow \end{array},$$

$$\langle n;0|H' \hat{\Pi}_n^1 H'|n;0\rangle \xrightarrow{(1.89c)} E_n^{(2)} = \begin{array}{c} n \quad m \quad n \\ \longrightarrow \otimes \longrightarrow \otimes \longrightarrow \end{array},$$

$$\langle n;0|H' \hat{\Pi}_n^1 H' \hat{\Pi}_n^1 H'|n;0\rangle - \langle n;0|H' \hat{\Pi}_n^2 H'|n;0\rangle \langle n;0|H'|n;0\rangle$$

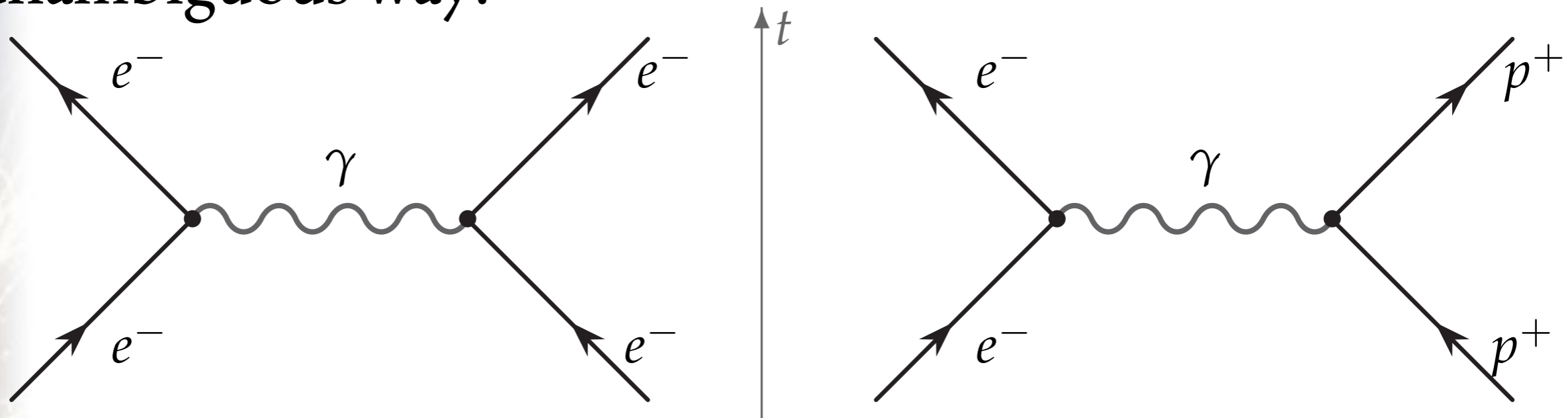
$$\xrightarrow{(1.89e)} E_n^{(3)} = \begin{array}{c} n \quad m \quad m \quad n \\ \longrightarrow \otimes \longrightarrow \otimes \longrightarrow \otimes \longrightarrow \end{array} - \begin{array}{c} n \quad n \\ \longrightarrow \otimes \longrightarrow \\ n \quad m \quad n \\ \longrightarrow \otimes \longrightarrow \otimes \longrightarrow \end{array},$$



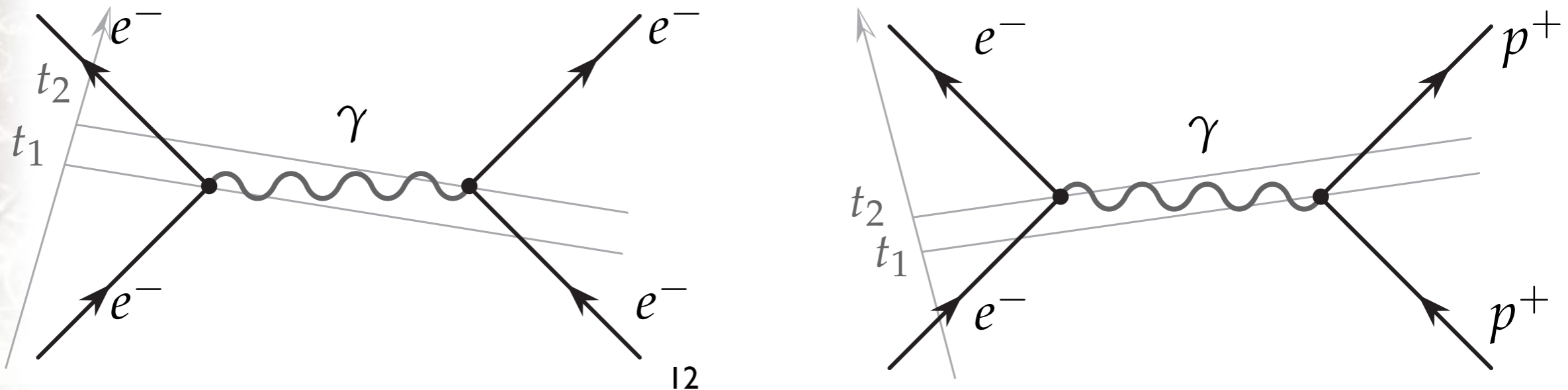
# Perturbation Picturebook

## SPACETIME PROCESSES (QM $\rightarrow$ QFT)

- The general idea is to depict physical processes, in a 1–1 unambiguous way:



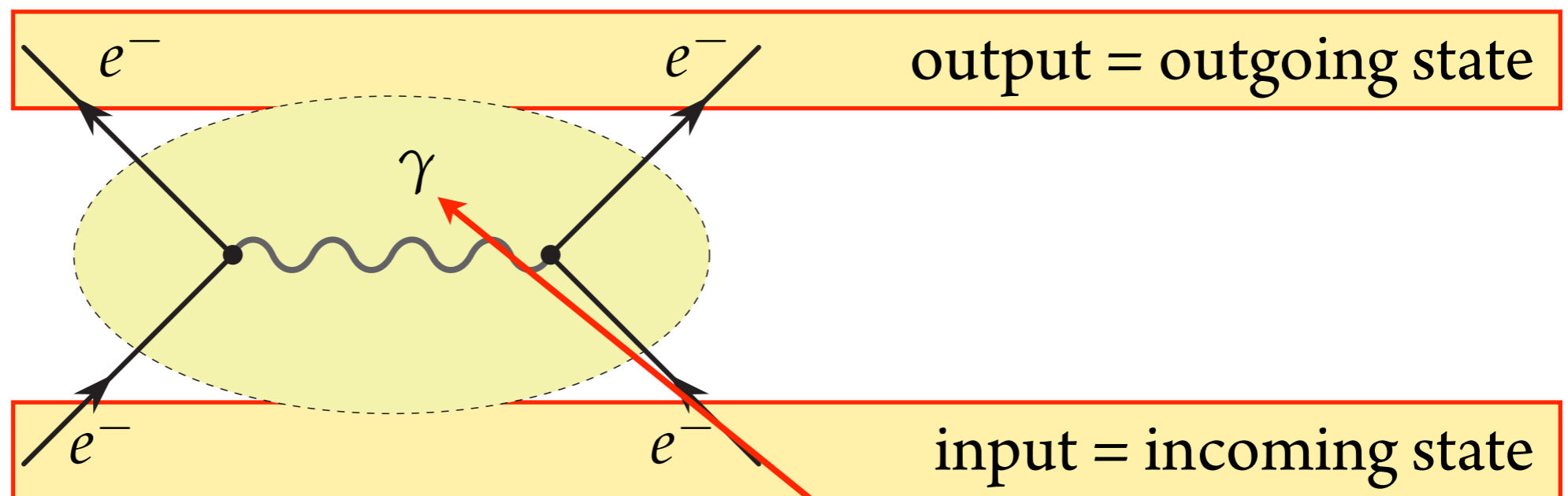
- ...but not literally!



# Perturbation Picturebook

## SPACETIME PROCESSES (QM $\rightarrow$ QFT)

- The “process” is whatever takes “input”  $\rightarrow$  “output”:

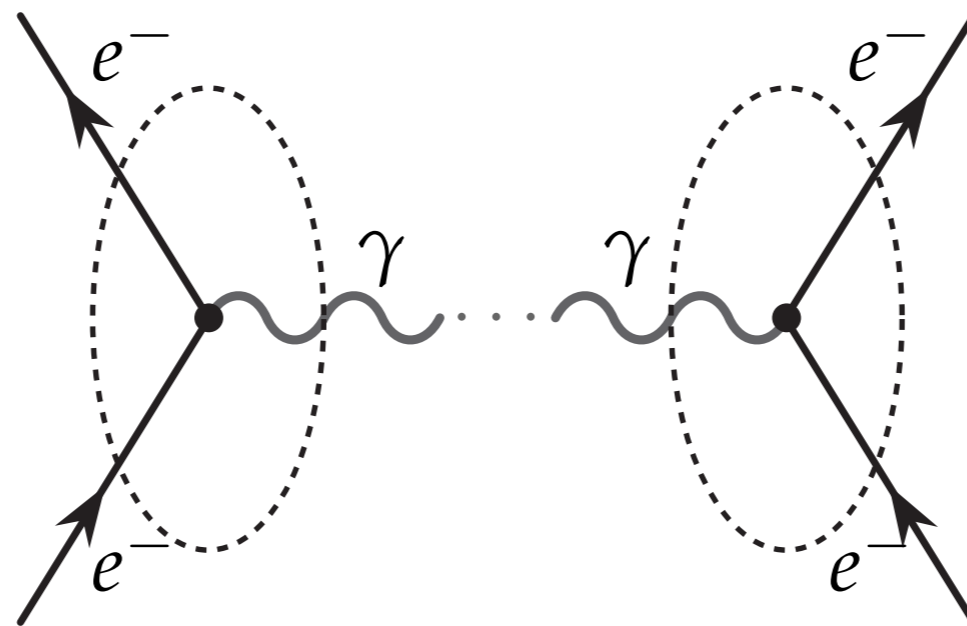


- In this case, it is a 1-photon electromagnetic interaction (scattering) of an electron with another electron.
- What is observed are the incoming two particles in the incoming 2-particle state, and the outgoing two particles in the outgoing 2-particle state. **Not the exchanged  $\gamma$ !**

# Perturbation Picturebook

## SPACETIME PROCESSES (QM $\rightarrow$ QFT)

- The un-observable exchange photon is thus virtual!
- Indeed, the two 3-particle processes separately...

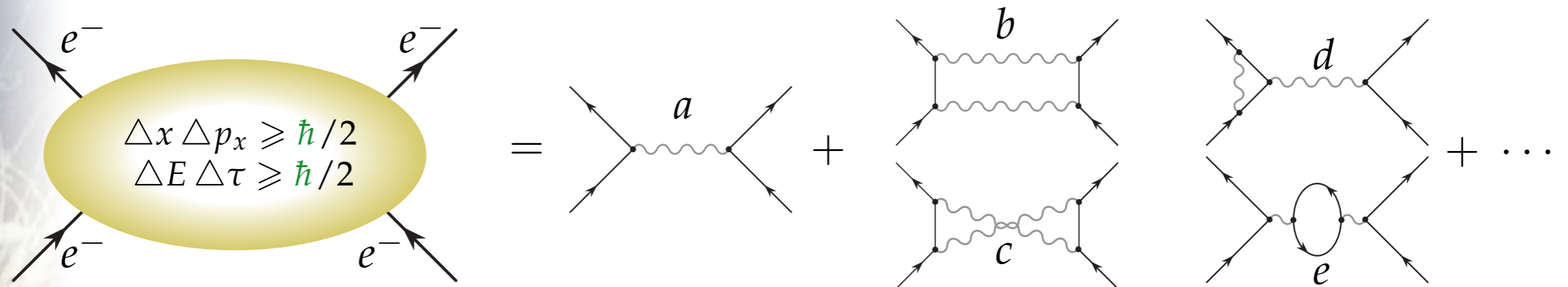


- ... are kinematically forbidden.
- Q.: Why?
- A.: Conservation of 4-momentum.

# Perturbation Picturebook

## SPACETIME PROCESSES (QM $\rightarrow$ QFT)


- The “process” is whatever takes “input”  $\rightarrow$  “output”:



- *Whatever is not forbidden, is mandatory. –R. Feynman*
- Notice, these sub-processes can be ordered:
  - by counting the “elementary” interaction vertices
  - by counting loops
  - by (non-)planarity
- ... which are qualitative—**& depicted**—characteristics

# Perturbation Picturebook

## SPACETIME PROCESSES (QFT)

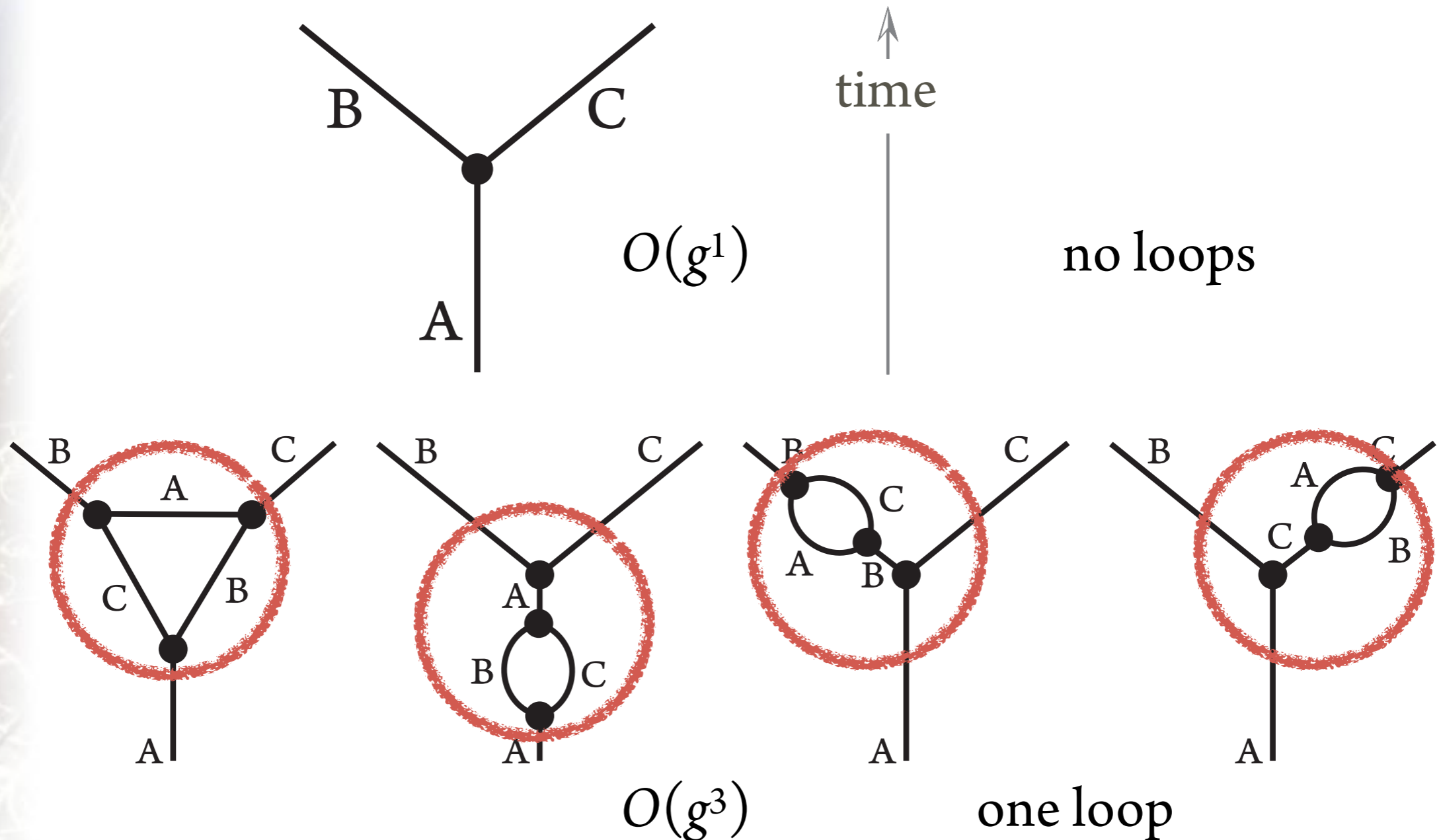
- A 1–1 correspondence:
  - the fundamental theory that designs the considered process,
  - diagram elements depicting terms from a Lagrangian,
  - rules of linking graphical elements, depicting a computation with the individual terms from the specified Lagrangian,
  - rules of listing all possible—and needed—Feynman diagrams,
  - the final mathematical expression for the matrix element (amplitude of probability) for the considered process, as a weighted sum of sub-processes,
  - the computation (estimate) of this mathematical expression.
- This would be the goal of a Quantum Field Theory course ...
- ... which this is not. 



# Perturbation Picturebook

## SPACETIME PROCESSES (QFT)

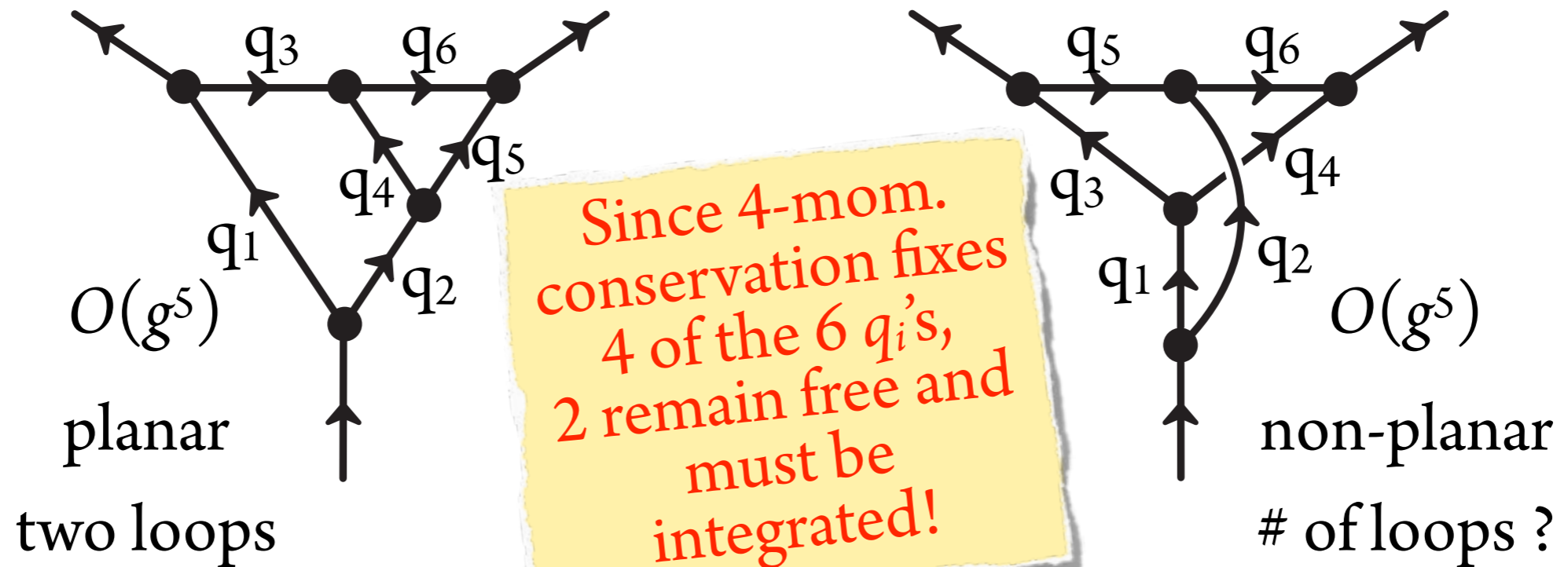
- Back in the  $A$ - $B$ - $C$  toy model, the  $A \rightarrow B+C$  decay.



# Perturbation Picturebook

## SPACETIME PROCESSES (QFT)

- And from  $O(g^5)$  onward, there are also complications:



$$\text{Solve}[\{z_1 = q_1 + q_2, q_1 = z_2 + q_3, q_2 = q_4 + q_5, q_4 + q_3 = q_6, q_5 + q_6 = z_3, z_1 = z_2 + z_3\}]$$

$$\{ \begin{aligned} q_2 &\rightarrow -q_1 + z_2 + z_3, \\ q_4 &\rightarrow -q_1 + q_6 + z_2, q_5 &\rightarrow -q_6 + z_3, \\ q_3 &\rightarrow q_1 - z_2, z_1 &\rightarrow z_2 + z_3 \end{aligned} \}$$

$$\text{Solve}[\{z_1 = q_1 + q_2, q_1 = q_3 + q_4, q_3 = z_2 + q_5, q_4 + q_6 = z_3, q_2 + q_5 = q_6, z_1 = z_2 + z_3\}]$$

$$\{ \begin{aligned} q_1 &\rightarrow q_5 - q_6 + z_2 + z_3, \\ q_3 &\rightarrow q_5 + z_2, q_4 &\rightarrow -q_6 + z_3, \\ q_2 &\rightarrow -q_5 + q_6, z_1 &\rightarrow z_2 + z_3 \end{aligned} \}$$

# A Few Stray Comments

## OVERVIEW

Molecular Physics	molecules (bound states of atoms) (foundation of all chemistry!)
Atomic Physics	atoms ( $e^-$ & nuclei bound states)
Nuclear Physics	atomic nuclei ( $p^+$ & $n^0$ bound states)
Elementary Particle Physics	hadrons (quark bound states)
	quarks, leptons, photons,...

Notice the extraordinary double-duty done by the physics discipline of Elementary Particle Physics—possibly including also a third tier, as professed to exist by such as string theory.

# Perturbation Picturebook

## CONSERVATION LAWS

- Strict Conservation Laws
  - Spacetime
    - Continuous (4-momentum & angular momentum)
    - Discrete ( $P, T, C$ )
  - Not Spacetime
    - EM charge
    - Chromodynamic Color
    - Weak Isospin
- Phenomenological Conservation Laws
  - lepton number(s, incl.  $e^-$ ,  $\mu^-$  &  $\tau^-$  separately)
  - baryon number(s, incl. strangeness, charm, ...)

# Thanks!

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