Introspections on Quantum Introspection

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(a)

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Quantum Introspections

Playbill

- Axioms and Assumptions
- Binary Breakdown?
- Critical Coaxing
- Division Drill
- Elusive Extensions
- Fusing Force-Fields
- Grand-Gödelian Guide



BRAND CAMP

by Tom Fishburne

THE FIVE STAGES OF RECESSION KEEP CALM HOW THE MY PLANS IT WILL BE AND CARRY. HECK AM I 1.50 A QUICK ARE 21 SUPPOSED TO ON RELOVERY RUINED 10 DO MORE WITH LESS? 2.12 1714 100 -4-BARGAINING ACCEPTANCE DENIAL ANGER DEPRESSION @ 2009 TOM FISHBURNE COM

"Whereof we cannot speak clearly, we must shut up." —w/apologies to L. Wittgenstein

Axioms and Assumptions

In Quantum Mechanics:

(1) Every state of a system \leftrightarrow a "state function," $|\psi\rangle \in \mathcal{S}$

(2) Every observable \leftrightarrow an operator, Q, acting on $|\psi\rangle$'s

(3) Only eigenvalues of Q are results of any individual measurement

(4) If $Q|n\rangle = q_n |n\rangle$, then $\operatorname{Prob}(Q \mapsto q_n |\psi) = \frac{|\langle n|\psi\rangle|^2}{\langle n|n\rangle\langle\psi|\psi\rangle} = \text{``}\cos^2(\theta_{n\psi})\text{''}$

(5) $i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$ with H the Hamiltonian operator (6) $|\psi\rangle \xrightarrow{Q \mapsto q_n} |n\rangle$

One more thing ([©]):

(7) The Hamiltonian, H, is independent of $|\psi\rangle$



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Axioms and Assumptions

• BTW, "(7) $H \neq H(\psi)$ " \Rightarrow "(5) $i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$ " a linear ODE

• \Rightarrow superposition, i.e., \mathcal{S} is a linear vector space

• Implicit: $\langle \psi | \chi \rangle < \infty$, by "(4) $\operatorname{Prob}(\mathbb{Q} \mapsto q_n | \psi) = \frac{|\langle n | \psi \rangle|^2}{\langle n | n \rangle \langle \psi | \psi \rangle}$,"

• So, \mathcal{S} is by (4) & for (7) a Hilbert space (...complete? ...rigged? ...)

• ... on which Q act as linear operators (... & bounded ...)

• (1) lied (): "mixed" states $\leftrightarrow |\psi\rangle$; $\rho := \sum_{i} r_i |i\rangle\langle i|$

Incomplete Most Q are Hermitian; & other, if eigenvalues are real

 $(5) \Rightarrow$ probability conservation; & particle decay? & particle condensation? & particle creation?

Binary Breakdown?

- Measurement conundrum:
 - How exactly does "(6) $|\psi\rangle \xrightarrow{Q \mapsto q_n} |n\rangle$ " happen?
 - What, who, when, where, why (&how) ... measures?
 - Collapse? Decoherence? Non-unitary? Non-local?
 - Parallel Universes? (combinatorial → transfinite?) *Mind? Friend?* ...?
 - "Be bothered sleepless, or have rocks in one's head" [~D. Mermin/TH]
 - The transition $(|\psi\rangle \rightarrow |n\rangle) \in \mathcal{S}$ is discontinuous...
 - Contradicts $i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$, i.e., $|\psi(t)\rangle = \exp\{-\frac{1}{\hbar}\int_0^t d\tau H\} |\psi(0)\rangle !!$
 - Of course, not H = H(t) controls the *t*-discontinuity





...getting our hands dirty...

Binary Breakdown?

- OK, $|\psi(t)\rangle = \exp\left\{-\frac{1}{\hbar}\int_{0}^{t} d\tau H(\tau)\right\} |\psi(0)\rangle$ <u>can</u> be discontinuous
 - "quantum jumps" including the "collapse" in measurements
 - as *modeled* by the *choice* of H(t) & by the choice of BC
- So, how *do* we choose/model H(t) ?
 - Typically, $H(t) = H_0 + V(t)$, $H_0 = KE + V_0$
 - Only " $KE = \frac{1}{2m} P^2$ " is *intrinsic* to the quantum system at hand
 - Both $V_0 \& V(t)$ encode interaction with its *environment*
 - Like "potential well," "Coulomb potential," ...
 -which is not quantal, but is treated as if classical

...getting our hands dirty...

Binary Breakdown?

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....getting our hands dirty....

Binary Breakdown?

- So, $|\psi(t)\rangle = \exp\{-\frac{1}{\hbar}\int_0^t d\tau H(\tau)\} |\psi(0)\rangle$ <u>can</u> have solutions
 - and <u>can</u> model the discontinuous change in $|\psi(t)\rangle$ under the influence...

…I'll be back!

- ... of a measuring device represented by the "external" V(t)
- How precisely *does* discontinuous change occur?
 - The "quantized" subsystem $\mapsto |\psi(t)\rangle$
 - The "classical" subsystem = environment/device/...
 - but ...just how do "quantum" and "classical" interact? (since they are inherently diametrically different)
 - Reminds of the unobservable "Führungsfeld" (pilot-wave)
 - that encodes all quantum interference & guides observable particles
 - by way of a never specified coupling/interaction to it (& which one?)

I have seen the problem. — It is us.

....getting our hands dirty....

Critical Coaxing

"Let's treat the environment also quantum-ly!"

- If the "object" $\mapsto |\psi(t)\rangle$, its "environment" $\mapsto |\phi(t)\rangle$
- <u>Toy model</u>: $i\hbar \frac{d}{dt} |\psi\rangle = V_M |\psi\rangle$, $V_M = (classical)$ measuring device
- ...and vice versa: $i\hbar \frac{d}{dt} |\phi\rangle = \Lambda_M |\phi\rangle$, $|\phi\rangle = state of the device$
- ... but then $V_M = V_M(\phi)$ and $\Lambda_M = \Lambda_M(\psi)$
- Coarsely: $|\psi\rangle \mapsto \psi(...,t)$

• so:
$$i\hbar \frac{\mathrm{d}\psi}{\mathrm{d}t} = (a_0 + a_1\phi + \dots)\psi$$

 $i\hbar \frac{\mathrm{d}\phi}{\mathrm{d}t} = (b_0 + b_1\psi + \dots)\phi$



Critical Coaxing

• This is a coupled system of *t*-ODEs

$$i\hbar \frac{\mathrm{d}\psi}{\mathrm{d}t} = (a_0 + a_1\phi + \dots)\psi^{-V_M}(\phi)$$
$$i\hbar \frac{\mathrm{d}\phi}{\mathrm{d}t} = (b_0 + b_1\psi + \dots)\phi^{-\Lambda_M}(\psi)$$



• Toy model (coarse): nonlinearly coupled object-measurer interaction

• So,
$$\ddot{\psi} = a_0 \left(\frac{b_0}{\hbar^2} + \frac{\dot{a}_0}{i\hbar a_0} - \frac{\dot{a}_1}{i\hbar a_1} \right) \psi + \left(\frac{\dot{a}_1}{a_1} + \frac{b_0}{i\hbar} \right) \dot{\psi}$$

 $+ \frac{a_0 b_1}{\hbar^2} \psi^2 + \frac{b_1}{i\hbar} \psi \dot{\psi} + \frac{\dot{\psi}^2}{\psi} + \dots$

• "QM is Either Non-Linear or Non-Introspective" [quant-ph:9712047]

STAGES OF GRIEF





"QM is either Non-Linear or Non-Introspective" [quant-ph:9712047]

→ Mod. Phys. Lett. A13 (1998) 2503–2512

...getting our hands really dirty...



Iteratively solving mutually interactive ("observing") system

... getting our hands really dirty...

Division Drill

• So, either

A nonlinearly coupled system of two equations

 $\int_{n} \underbrace{\left[-\frac{\hbar^{2}}{2m_{e}}\overrightarrow{\nabla}^{2}-\frac{e^{2}}{4\pi\epsilon_{0}}\left(\frac{1}{r}+\frac{1}{\sqrt{r^{2}+R^{2}-2rR\cos\theta}}\right)\right]}{\left[-\frac{\hbar^{2}}{2m_{p}}\overrightarrow{\nabla}^{2}_{R}+\frac{2e^{2}}{4\pi\epsilon_{0}}\frac{1}{|\overrightarrow{R}|}+V_{e}(\overrightarrow{R})\right]}\phi(\overrightarrow{R})=E_{p}\phi(\overrightarrow{R})$

solved iteratively

• Or, a "unified" description: $-\left[\frac{\hbar^{2}}{2m_{e}}\overrightarrow{\nabla_{e}}^{2} + \frac{\hbar^{2}}{2m_{p}}(\overrightarrow{\nabla_{a}}^{2} + \overrightarrow{\nabla_{b}}^{2}) + \frac{Escher's view}{2m_{e}}\right] (Escher's view)$ $+\frac{e^{2}}{4\pi\epsilon_{0}}\left(\frac{1}{|\overrightarrow{r_{e}} - \overrightarrow{r_{a}}|} + \frac{1}{|\overrightarrow{r_{e}} - \overrightarrow{r_{b}}|} - \frac{1}{|\overrightarrow{r_{a}} - \overrightarrow{r_{b}}|}\right) \Psi(\overrightarrow{r_{e}}, \overrightarrow{r_{a}}, \overrightarrow{r_{b}}) = E_{M}\Psi(\overrightarrow{r_{e}}, \overrightarrow{r_{a}}, \overrightarrow{r_{b}})$

Vishut up and calculate."

Insider View

"...nobody really understands quantum mechanics." -R.P. Feynman

Who, what, when, where ... is doing the measuring? *How & why does superposition ... break?*

Elusive Extensions

Bohr, Heisenberg, Wigner, Bohr, DeWitt, • Observation irreversibly collapses (*dis*-superposes) the state • The observometer? The human viewer? Her friend? Twice removed? • Objective-collapse by ... nonlinearities and/or mesoscale dynamics • Gravity? Spontaneous localization? Origin of the random noise? Environmental decoherence by statistically many scattering events • Quantum Prob \rightarrow Classical Prob, but still linear QM & superposition? Incessant splitting of (increasingly) many worlds • Coexistent with decoherence, but indistinguishable from Copenhagen? Angelo Bassi et al.: "Models of Wave-function Collapse, Underlying Theories, and Experimental Tests" Rev. Mod. Phys. 85 (2013) 471-527, arXiv:1204.4325

Bohr, Heisenberg, Wigner, Bohr, DeWitt, Who, what, when, where ... is doing the measuring? *How & why does superposition ... break?*

Elusive Extensions

- Modifications of QM ... often include nonlinearities
 - Tend to be superluminal? Stochastic prob. distribution *ad hoc*?
- Hidden variables... Trace Dynamics (Themodynamical QM)...
 - Nonlocality? Superluminal? Grassmann matrices? (\rightarrow M-Theory?)
- Spacetime: emergent? smeared? non-commutative?
 - Heisenberg? (incl. "Loop Quantum Gravity"?)
- Unseen Führungsfeld (pilot-wave) guides observable particles
 - By what interaction? What mediates that? Is that particle-specific? (How do the various Führungfeldern interact? ... & avoid other particles?)



"Quantum Field Theory already is what Führungsfeldtheorie wants to become when it grows up." (Pilot-wave theory)

originally so-named by Max Born in 1926



- 1972 S. Freedman & J. Clauser's 1st experiment violating Bell('64)
 Clauser-Horne-Shimony-Holt ('69) inequalities; A. Aspect('81)...
- 1974: The Standard Model *quantum consistency* 1969 Adler-Bell-Jackiw & 1970 Glashow-Iliopoulos-Maiani
 - 1974, Ting & Richter: $J/\psi = [c\bar{c}]$; "November revolution"
 - A concrete QFT model involving
 - Renormalizable non-abelian quantum gauge theory fundamental forces
 - ...with spontaneous symmetry breaking
 - ...experimentally detected 2011–13 (Higgs boson)
 •••I'll be back!



 \rightarrow charm

Standard Model of FUNDAMENTAL PARTICLES AND INTERACTIONS

The Standard Model summarizes the current knowledge in Particle Physics. It is the quantum theory that includes the theory of strong interactions (quantum chromodynamics or QCD) and the unified theory of weak and electromagnetic interactions (electroweak). Gravity is included on this chart because it is one of the fundamental interactions even though not part of the "Standard Model."

matter constituents FERMIONS spin = 1/2, 3/2, 5/2, ...

Leptons spin = 1/2				Quarks spin = 1/2			
Flavor	Mass GeV/c ²	Electric charge	ctric Flavor		Approx. El Mass El GeV/c ² ch		
ve electron neutrino	<1×10 ⁻⁸	0		U up	0.003	2/3	
e electron	0.000511	-1		d down	0.006	-1/3	
ν_{μ} muon neutrino	<0.0002	0		C charm	1.3	2/3	
μ muon	0.106	-1		S strange	0.1	-1/3	
v_{τ} tau neutrino	<0.02	0		t top	175	2/3	
τ tau	1.7771	-1		b bottom	4.3	-1/3	

Spin is the intrinsic angular momentum of particles. Spin is given in units of \hbar , which is the quantum unit of angular momentum, where $\hbar = h/2\pi = 6.58 \times 10^{-25}$ GeV s = 1.05×10^{-34} J s.

Electric charges are given in units of the proton's charge. In SI units the electric charge of the proton is 1.60×10^{-19} coulombs.

The energy unit of particle physics is the electronvolt (eV), the energy gained by one electron in crossing a potential difference of one volt. Masses are given in GeV/ c^2 (remember $E = mc^2$), where 1 GeV = 10⁹ eV = 1.60×10⁻¹⁰ joule. The mass of the proton is 0.938 GeV/ c^2 = 1.67×10-27 kg.

Baryons qqq and Antibaryons qqq Baryons are fermionic hadrons. There are about 120 types of baryons.						
Symbol	Name	Quark content	Electric charge	Mass GeV/c ²	Spin	
р	proton	uud	1	0.938	1/2	
p	anti- proton	ūūd	-1	0.938	1/2	
n	neutron	udd	0	0.940	1/2	
Λ	lambda	uds	0	1.116	1/2	
Ω-	omega	555	-1	1.672	3/2	

Matter and Antimatter

For every particle type there is a corresponding antiparticle type, denoted by a bar over the particle symbol (unless + or - charge is shown). Particle and antiparticle have identical mass and spin but opposite charges. Some electrically neutral bosons (e.g., Z^0 , γ , and $\eta_c = c\bar{c}$, but not $K^0 = ds$) are their own antiparticles.

Figures

These diagrams are an artist's conception of physical processes. They are not exact and have no meaningful scale. Green shaded areas represent the cloud of gluons or the gluon field, and red lines the quark paths.



PROPERTIES OF THE INTERACTIONS

BOSONS

Unified Electroweak spin = 1				
Name	Mass GeV/c ²	Electric charge		
γ photon	0	0		
W-	80.4	-1		
W+	80.4	+1		
Z ⁰	91,187	0		

force carriers spin = 0, 1, 2, ...

Ele	ctroweak	spin = 1	Stro
	Mass GeV/c ²	Electric charge	Name
n	0	0	g gluon
	80.4	-1	Color Cha
	80.4	+1	Each quark "strong char These chara
	80.4 80.4	-1 +1	Color Ch Each quark "strong ch These char



rge

carries one of three types of ge," also called "color charge." have nothing to do with the colors of visible light. There are eight possible types of color charge for gluons. Just as electri-

cally-charged particles interact by exchanging photons, in strong interactions color-charged particles interact by exchanging gluons. Leptons, photons, and **W** and **Z** bosons have no strong interactions and hence no color charge.

Quarks Confined in Mesons and Baryons

One cannot isolate quarks and gluons; they are confined in color-neutral particles called hadrons. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional guark-antiguark pairs (see figure below). The guarks and antiguarks then combine into hadrons; these are the particles seen to emerge. Two types of hadrons have been observed in nature: mesons gg and baryons ggg.

Residual Strong Interaction

The strong binding of color-neutral protons and neutrons to form nuclei is due to residual strong interactions between their color-charged constituents. It is similar to the residual electrical interaction that binds electrically neutral atoms to form molecules. It can also be viewed as the exchange of mesons between the hadrons.

							M	
Property	Gravitational	Weak Electromagnetic (Electroweak)		Strong		Mesons / There are ab		
				Fundamental	Residual			
Acts on:	Mass – Energy	Flavor	Electric Charge	Color Charge	See Residual Strong Interaction Note	Symbol	Name	
Particles experiencing:	All	Quarks, Leptons	Electrically charged	Quarks, Gluons	Hadrons	77+	nion	
Particles mediating:	Graviton (not yet observed)	W+ W- Z ⁰	γ	Gluons	Mesons	W-	pion	
Strength relative to electromag 10 ⁻¹⁸ m	10-41	0.8	1	25	Not applicable	•	kaon	3
for two u quarks at: 3×10 ⁻¹⁷ m	10-41	10-4	1	60	to quarks	ρ^+	rho	U
ter two protons in nucleus	10-36	10-7	1	Not applicable	20	B0	B-zero	c

Mesons qq Mesons are bosonic hadrons. There are about 140 types of mesons.							
Symbol	Symbol Hame Quark Electric Mass GeV/c ² Spir						
π^+	pion	uđ	+1	0.140	0		
К-	kaon	sū	-1	0.494	0		
ρ^+	rho	uđ	+1	0.770	1		
B ⁰	B-zero	db	0	5.279	0		
$\eta_{\rm c}$	eta-c	cī	0	2 .980	0		

e+e- -> B0 B0 n→pe⁻ v_o e⁻ e An electron and positron antielectron) colliding at high energy can nnihilate to produce B⁰ and B⁰ mesons A neutron decays to a proton, an electron and an antineutrino via a virtual (mediating) W boson. This is neutron ß decay. ia a virtual Z boson or a virtual photor



Two protons colliding at high energy can produce various hadrons plus very high mass particles such as Z bosons. Events such as this one are rare but can yield vital clues to the structure of matter

hadrons

to hadrons

hadrons

Z⁰

The Particle Adventure Visit the award-winning web feature The Particle Adventure at http://ParticleAdventure.org

This chart has been made possible by the generous support of:

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"Quantum Field Theory already is what Führungsfeldtheorie wants to become when it grows up."

Fusing Force-Fields

- Paraphrasing "The pilot-wave interpretation of quantum mechanics is wishful thinking that aims to achieve what quantum field theory has been for the past half a century." [Quora answers: 1 & 2]
- Quantum fields exist continuously throughout spacetime
- Change in the field is quantized \rightarrow [particle, ...packet,... wave]
- *Fermions*: Pauli-exclude from condensing \rightarrow "background Fermi field" = \emptyset
- *Bosons*: can condense (in any particular one state) \rightarrow "background Bose field"
- Very large ensemble ("sea") of quanta comprises the continuous field
- A field and its quanta do not *interact* with each other...
 - ... they are part and parcel of *the same entity*

(no need for velcro)

- I974: (Super)string theory = theory of gravity
 - A finite inherently quantum theory
 - with gauge fields and fermions and (quantum!) gravity
 - 1979 (D. Friedan): *quantum stability* \Rightarrow Einstein equations
 - 1984 (Alvarez-Gaumé & Witten): gravitational anomalies & cancellation
 - 1984 (Green-Schwarz): one more *anomaly cancellation*
 - 1984 (GHMR): heterotic string model with $E_8 \times E_8$ gauge fields
 - 1984 (CHSW): supersymmetric stability \Rightarrow Calabi-Yau compactification
 - 1984 me@UMD 1987: $O(10^3)$ models; \rightarrow UT: '88 connected web connected by phase transitions



quantum consistency

- Quantum Field Theory
 - A <u>framework</u> to construct concrete <u>models</u>
 - The Standard Model has ~26 continuous parameters & many other choices

- What is it that *is being* observed/measured?
 - Such as in the 2011–13 discovery of the Higgs boson?





- Quantum Field Theory
 - A *framework* to construct concrete *models*
 - The Standard Model has ~26 continuous parameters & many other choices
 - The QFT framework is axiomatic (Wightman; Osterwalder-Schrader;
- What is it that *is being* observed/measured?
 - Such as in the 2011–13 discovery of the Higgs boson?
 - Terabytes upon terabytes of data from hierarchically triggered and computer-controlled layers of tons of detectors ...stored until data-mined
 - Relative process amplitudes: common factors cancel (~Wigner-Eckart) and are largely ...irrelevant





Haag-Kastler; ...)



The Big Picture

Grand-Gödelian Guide

• QM: 6 (oft-quoted) axioms & (7) $H \neq H(\psi)$ in $i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$

- (7) \Rightarrow linear superposition \Rightarrow quantum entanglement/non-factorizability
- But the division drill \Rightarrow nonlinearity \Rightarrow no superposition
- which is then, logically, the "axiom (¬7)"
 The six axioms that is the image of the six axioms of the six axioms
- The six axioms thereby imply neither "axiom (7)" nor "axiom $(\neg 7)$ "
 - (1) Every state of a system \leftrightarrow a "state function," $|\psi\rangle \in \mathcal{S}$
 - (2) Every observable \leftrightarrow an operator, Q, acting on $|\psi\rangle$'s
 - (3) Only eigenvalues of Q are results of any individual measurement
 - (4) If $Q | n \rangle = q_n | n \rangle$, then $\operatorname{Prob}(Q \mapsto q_n | \psi) = \frac{|\langle n | \psi \rangle|^2}{\langle n | n \rangle \langle \psi | \psi \rangle} = \text{``} \cos^2(\theta_{n\psi})$ '' (5) $i\hbar \frac{\mathrm{d}}{\mathrm{d}t} | \psi \rangle = \mathsf{H} | \psi \rangle$ with H the Hamiltonian operator

 - (6) $\Pi(\mathbf{Q} \mapsto q_n) : |\psi\rangle \to |n\rangle$

... getting our hands really dirty ..

Division Drill

• So, either



• A nonlinearly coupled system of two equations $-\left[\frac{\hbar^2}{2m_e}\vec{\nabla}^2 + \frac{e^2}{4\pi\epsilon_0}\left(\frac{1}{r} + \frac{1}{\sqrt{r^2 + R^2 - 2rR\cos\theta}}\right)\right]\psi(\vec{r}) = E_e\psi(\vec{r}) + 1$ $-\left[\frac{\hbar^2}{2m_p}\vec{\nabla}_R^2 + \frac{2e^2}{4\pi\epsilon_0}\frac{1}{|\vec{R}|} + V_e(\vec{R})\right]\phi(\vec{R}) = E_p\phi(\vec{R}) + 1$

solved iteratively

• Or, a combined description: $-\left[\frac{\hbar^{2}}{2m_{e}}\overrightarrow{\nabla}_{e}^{2} + \frac{\hbar^{2}}{2m_{p}}(\overrightarrow{\nabla}_{a}^{2} + \overrightarrow{\nabla}_{b}^{2}) \quad (Escher's view) + \frac{e^{2}}{4\pi\epsilon_{0}}\left(\frac{1}{|\overrightarrow{r}_{e} - \overrightarrow{r}_{a}|} + \frac{1}{|\overrightarrow{r}_{e} - \overrightarrow{r}_{b}|} - \frac{1}{|\overrightarrow{r}_{a} - \overrightarrow{r}_{b}|}\right)\right]\Psi(\overrightarrow{r}_{e}, \overrightarrow{r}_{a}, \overrightarrow{r}_{b}) = E_{t}\Psi(\overrightarrow{r}_{e}, \overrightarrow{r}_{a}, \overrightarrow{r}_{b})$



Grand-Gödelian Guide

- In fact Gödelian incompletenes
 = <u>extensibility</u> ∈ Physics
 - System of equations = a particular model in classical fluid mechanics
 - \rightarrow Maxwell's equations
- Just as a QFT model with only two families of fermions
 - or a strong $SO(3)_g$ -force (Higgs-broken from $SU(3)_c$) *Phys. Rev.* D33 (1986) 1429 erratum: *Phys. Rev.* D34 (1986) 3536

$e + \frac{df}{dt} + \frac{dg}{dt} + \frac{dh}{dt} = 0$	(1)	Gauss' Law
dx dy dz	- 21	
$\mu \alpha = \frac{dH}{dy} - \frac{dG}{dz}$ $\mu \beta = \frac{dF}{dz} - \frac{dH}{dx}$ $\mu \gamma = \frac{dG}{dx} - \frac{dF}{dy}$	(2)	Equivalent to Gauss' Law for magnetism
$P = \mu \left(\gamma \frac{dy}{dt} - \beta \frac{dz}{dt} \right) - \frac{dF}{dt} - \frac{d\Psi}{dz}$ $Q = \mu \left(\alpha \frac{dz}{dt} - \gamma \frac{dx}{dt} \right) - \frac{dG}{dt} - \frac{d\Psi}{dy}$ $R = \mu \left(\beta \frac{dx}{dt} - \alpha \frac{dy}{dt} \right) - \frac{dH}{dt} - \frac{d\Psi}{dz}$	(3)	Faraday's Law (with the Lorentz Force and Poisson's Law)
$\frac{d\gamma}{dy} - \frac{d\beta}{dz} = 4\pi p' \qquad p' = p + \frac{df}{dt}$ $\frac{d\alpha}{dz} - \frac{d\gamma}{dx} = 4\pi q' \qquad q' = q + \frac{dg}{dt}$ $\frac{d\beta}{dx} - \frac{d\alpha}{dy} = 4\pi r' \qquad r' = r + \frac{dh}{dt}$	(4)	Ampère-Maxwell Law
$\mathbf{P} = -\xi p \mathbf{Q} = -\xi q \mathbf{R} = -\phi$		Ohm's Law
P = kf Q = kg R = kh		The electric elasticity equation ($\mathbf{E} = \mathbf{D}/\epsilon$)
$\frac{de}{dt} + \frac{dp}{dx} + \frac{dq}{dy} + \frac{dr}{dz} = 0$	- 1	Continuity of charge

Grand-Gödelian Guide

• How to construct *Models* in *Theories*?

- The domain space, \mathfrak{D} ,
- The target space, $\mathfrak{T}, \leftarrow presumed!$
- The map $\varphi \colon \mathfrak{D} \to \mathfrak{T}$
- The functional $S[\varphi; C] := \int_{\mathfrak{N}} \mathscr{L}(\varphi, \dot{\varphi}, ...; C) \& BC$
- Classical theory: $\varphi_{cl} =$ solution of $\delta_{\varphi}S[\varphi; C] = 0$
 - Quantum theory: (sum over all alternative histories) • The partition functional, $Z[\vartheta; C] := \int \mathbf{D}[\varphi] e^{(S[\varphi; C] + \int \vartheta \cdot \varphi)/i\hbar}$ • Correlations: $G(\xi_1, ..., \xi_n) := \lim_{\vartheta \to 0} \left\{ \frac{1}{Z[\vartheta; C]} \left[\frac{\partial}{\partial \vartheta(\xi_1)} \cdots \frac{\partial}{\partial \vartheta(\xi_n)} Z[\vartheta; C] \right] \right\}$

 $\dots etc$

 $\mathfrak{D} = \mathbb{R}_t^1; \quad \mathfrak{T} = \mathbb{R}_{\vec{r}}^3; \quad \varphi = \vec{r}(t);$ $S[\vec{r}; m, k] = \int dt \left\{ \frac{m}{2} \dot{\vec{r}}^2 - \frac{k}{2} \vec{r}^2 \right\}$

 $\mathfrak{D} = \mathbb{R}^{1,3}_{t,\vec{r}}; \quad \mathfrak{T} = \mathbb{R}^{1,3};$ $\varphi = (\Phi(\vec{r},t), \overrightarrow{A}(\vec{r},t))$

$S[\varphi; C] := \int_{\mathfrak{D}} \mathscr{L}(\varphi, \dot{\varphi}, ...; C) \& BC$ $Z[\vartheta; C] := \int \mathbf{D}[\varphi] e^{(S[\varphi; C] + \int \vartheta \cdot \varphi)/i\hbar}$ Grand-Gödelian Guide• We <u>can</u> turn the tables:

- Choose: \mathfrak{D} ; φ as "fields" on \mathfrak{D} ; $S[\varphi; C]$ by symmetries &c.
- Compute \mathfrak{T} as the range of so constrained φ
- Induce "layers" of models
 - Define $e^{\left(S_{\text{eff}}[\phi;\tilde{C}]\right)/i\hbar} := \int \mathbf{D}[\varphi] e^{\left(S[\varphi+\phi;C]\right)/i\hbar} \& \text{seek } \delta_Q C := (\tilde{C}-C) \stackrel{!}{=} 0$
 - 1979, Friedan: for $S[X^{\mu}; g_{\mu\nu}] = \int d^2\xi \gamma^{ij}(\xi) (\partial_i X^{\mu}) (\partial_j X^{\nu}) g_{\mu\nu}(X)$
 - $\delta_Q g_{\mu\nu}(X) \stackrel{!}{=} 0 \Rightarrow$ Einstein field equations!
 - 2nd layer
 ...which makes the Einstein-Hilbert action stationary
 - Generalized to all "familiar" gauge QFT models

Gödelian

worldsheet

spacet11

extensibility





The Biger Picture

Go Forth &

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